

Factors controlling contrail cirrus optical depth

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Abstract. Aircraft contrails develop into contrail cirrus by depositional growth and sedimentation of ice particles and horizontal spreading due to wind shear. Factors controlling this development include temperature, ice supersaturation, thickness of ice-supersaturated layers, and vertical gradients in the horizontal wind field. An analytical microphysical cloud model is presented and validated that captures these processes. Many individual contrail cirrus are simulated that develop differently owing to the variability in the controlling factors, resulting in large samples of cloud properties that are statistically analyzed. Contrail cirrus development is studied over the first four hours past formation, similar to the ages of contrails that were tracked in satellite imagery on regional scales. On these time scales, contrail cirrus optical depth and microphysical variables exhibit a marked variability, expressed in terms of broad and skewed probability distribution functions. Typical simulated mean optical depths at a wavelength of $0.55\text{ }\mu\text{m}$ are in the range $0.2 - 0.3$. A substantial fraction $20 - 40\%$ of contrail cirrus stay subvisible (optical depth < 0.02). A detailed analysis suggests that previous satellite measurements of line-shaped persistent contrails have missed about 86% (35%) of contrails with optical depth ≤ 0.05 ($0.05 - 0.1$), amounting to almost 50% of contrails of all optical depths. When comparing observations with simulations and when estimating the contrail cirrus climate impact, not only mean values but also the variability in optical depth and microphysical properties need to be considered.

1 Introduction

Knowledge of cloud optical depth is of paramount importance in any attempt to assess cloud-radiation interactions.

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Cloud types can be characterized by a range of altitudes and optical depths, e.g. according to the classification of the International Satellite Cloud Climatology Project (Rossow and Schiffer, 1999). Thin cirrus, a class of high clouds, are ubiquitous and can assume a broad range of low optical depth values. Thin cirrus with optical depths as low as 10^{-4} were observed by ground-based Lidar (e.g. Goldfarb et al., 2001) and have been routinely detected by space-borne sensors that measure the sunlight through the limb of the atmosphere (Wang et al., 1996). However, many thin cirrus remain undetected when using passive remote sensing methods that rely on upwelling radiance measurements (Wylie et al., 1995). Thin cirrus are important players in regulating the energy budget in the Earth-atmosphere system due mainly to trapping of longwave radiation, exerting a significant influence on the general circulation (e.g. Liou, 1986). These clouds may be present without visible manifestation and contain unfrozen aerosol particles besides ice condensate in the μm size range, emphasizing the continuum aspect of the aerosol/cloud system (Kärcher and Solomon, 1999).

The optical depth is key for determining shortwave cloud reflectance and absorption, longwave emissivities, and diabatic heating rates (e.g. Stephens, 1978). Cirrus radiative forcing is known to depend on horizontal inhomogeneity on scales $< 100\text{ km}$ (Fu et al., 2000; Carlin et al., 2002). Including inhomogeneity effects in global model simulations may substantially change radiative forcing (Gu and Liou, 2006). The latter is also sensitive to vertical inhomogeneity. Vertical variability is expressed in terms of variations in ice water content and effective ice crystal radii on scales $\sim 100\text{ m}$ (Miloshevich et al., 1997; Field and Heymsfield, 2003). However, very few studies have investigated the relationship between the spatial variability in optical depth and the underlying cloud physical processes at the process level (e.g. Hogan and Kew, 2005; Kay et al., 2006).

Aviation affects cirrus cloudiness by producing persistent contrails. Contrail cirrus originating from line-shaped con-

trails and developing in ice-supersaturated air belong to the category of thin cirrus. Line-shaped contrail coverage can be substantial on regional scales (e.g. Minnis, 2003) and individual contrails can exert radiative heating rates of 5–10 K/d (Jensen et al., 1998). In young contrails and possibly in contrail cirrus, the existence of numerous μm -sized ice crystals (aerosol/ice continuum) poses a great challenge for accurate in situ measurements. The difficulty of tracking contrails, once they have lost their linear shape, by aircraft or with satellite detection algorithms is another reason for the lack of robust data on contrail microphysical properties, and the apparent absence of measurements of older (age > 1 h) contrail cirrus.

Similar to natural cirrus, key factors controlling contrail cirrus development over many hours of lifetime include temperature, supersaturation, vertical shear of the horizontal wind, and the geometric dimensions of ice-supersaturated layers (Kärcher and Spichtinger, 2009). Radiative cooling or heating may affect the evolution of the ice phase and the moisture field primarily in optically thicker ice clouds (Dobbie and Jonas, 2001). While some effort has been made simulating contrail formation and early contrail evolution covering the first few minutes after formation, contrail cirrus evolution has not been systematically studied using process-based models. Such case studies may be based on high-resolution Large-Eddy simulations (LES). Interactions between contrail cirrus and the large-scale flow can only be investigated by means of large-scale models, in particular when the global climate impact of aviation-induced cloudiness is to be assessed (Burkhardt and Kärcher, 2009).

In this study, we investigate how contrail cirrus microphysical and optical properties depend on cloud-controlling factors and their atmospheric variability. We have performed systematic studies of contrail cirrus development under well defined atmospheric boundary conditions. Our study yields valuable statistical information on contrail cirrus properties. We underscore that such information can neither practically be obtained on a purely observational basis (because older contrail cirrus cannot easily be observed and the number of observations providing representative data would be prohibitive) nor by using cloud-resolving simulations (because the LES simulation of even a single cloud tracked over several hours requires large computational resources).

We seek a simplified analytical treatment of contrail cirrus that still captures the main physical mechanisms governing the contrail cirrus development (Sect. 2). This approach allows us to carry out studies of many thousands of individual contrail cirrus events in variable ambient conditions and to obtain large samples of cloud properties. We focus on statistics of optical depth but also discuss the underlying cloud properties, yielding important insight into the mean magnitude and variability of contrail cirrus parameters (Sect. 3). A comparison with observations is carried out, providing a better understanding of measurements of contrail properties (Sect. 4). An observational case study of statistics of con-

trail optical depth is analyzed in detail, showing that satellite measurements based on passive remote sensing strongly underestimate the occurrence frequency of optically thin contrails (Sect. 5). Section 6 presents a summary of our study, highlights our main results and points to areas of uncertainties to be addressed in future research.

2 Model and methods

The aim of this section is twofold. In Sect. 2.1, we discuss the simplified cloud model employed to estimate contrail cirrus evolution. This model is based on the dynamic equation describing the temporal and spatial evolution of the ice particle size distribution function and idealized assumptions concerning the meteorological fields in which the clouds evolve. The major physical processes considered are depositional growth and sedimentation of ice particles and their transport in a vertically sheared flow field. The dynamic equation is solved analytically resulting in sedimentation and growth histories of particles of different sizes. A representative explicit solution is compared to more detailed LES results over a range of contrail ages in order to judge the quality and usefulness of our model for contrail cirrus studies.

In Sect. 2.2, we introduce and discuss the probability distributions of the meteorological and aircraft-related factors controlling contrail cirrus development. This information is employed in Sect. 3 to drive a large number of individual analytical contrail cirrus simulations. The chosen functional forms of the distributions are largely based on observations and are intended to probe a wide range of environmental states affecting ice particle growth, sedimentation, and spreading. Further details concerning the model simulations including definitions of probability distributions of contrail cirrus properties are also given in Sect. 2.2.

2.1 Governing equation and basic assumptions

The dynamic evolution of the ice crystal mass distribution function $f(\mathbf{x}, m, t)$ is governed by

$$\partial_t f + \partial_m(\dot{m}f) + \nabla[(\mathbf{v} - \mathbf{v}_t)f] = \mathcal{T}[f], \quad (1)$$

where t denotes the time; \mathbf{x} and \mathbf{v} are the position and mean flow field velocity vectors with the horizontal and vertical components $\{x, y, z\}$ and $\{u, v, w\}$, respectively; \mathcal{T} is an operator summarizing effects of diffusion and turbulence on the particle distribution; and m , \dot{m} and \mathbf{v}_t are the ice crystal mass, mass growth rate and terminal fall speed, respectively. The second term on the left hand side describes ice crystal growth by diffusion of water vapor; growth by aggregation is considered less important, as most ice crystals in contrail cirrus stay relatively small (below $\sim 50 \mu\text{m}$) over the time scales considered here (Duda et al., 2001). The third term on the left hand side represents advective transport and sedimentation of ice crystals.

For a comprehensive solution, Eq. (1) would need to be coupled to equations governing the evolution of the flow and radiation fields, air temperature T and pressure p , and the ice saturation ratio S , or ice supersaturation $s = S - 1$, in the framework of a cloud-resolving model. We derive a simplified solution of Eq. (1) by introducing the following assumption. A vertically sheared horizontal flow field \mathbf{v} and constant homogeneous fields of T , p , and s are prescribed. Effects of turbulent diffusion on the ice particle size distribution and number concentration will be considered empirically in the solution of Eq. (1), which we obtain without explicit consideration of the operator \mathcal{T} . Because vertical air motion is not allowed to affect s , potential microphysical feedbacks, for instance caused by sedimentation of ice particles into clear air with rising relative humidity, are not considered.

Two approximations further simplify the problem. (i) The flow field is two-dimensional (2-D) with $\mathbf{x} = \{x, z\}$ and $\mathbf{v} = \{u, w\}$, allowing the horizontal velocity component $u(z)$ to vary with altitude in order to capture the effect of shear-induced contrail spreading, and with the terminal fall speed vector $\mathbf{v}_t = \{0, v_t\}$. (ii) Ice crystals are spherical, their mass is replaced by the particle radius r and the mass growth rate by a radial growth rate \dot{r} . The function f then represents the size distribution of ice crystals, i.e. the number of ice crystals per unit volume of air per unit particle radius.

2.1.1 Analytical solution

Depending on cloud age and location within cloud, natural cirrus may develop small ice particle modes at sizes $\lesssim 100 \mu\text{m}$ that are generated by ice nucleation and depositional growth (e.g. Schröder et al., 2000; Gayet et al., 2002) and large modes at sizes $\gtrsim 100 \mu\text{m}$ that consist of complex ice crystals formed by aggregation (e.g. Mitchell, 1991; Field and Heymsfield, 2003). Gamma functions have been proposed to represent the observed size distributions of small and large cirrus ice particle modes (Kosarev and Mazin, 1991) and have been frequently used in the literature ever since (e.g. Ivanova et al., 2001; Mitchell, 2002; Heymsfield, 2003; Delanoë et al., 2005; Morrison and Gettelman, 2008). Sampling of a sufficient amount of cloud elements in different cloud regions is required to construct such size distributions. In situ observations of aged (> 30 min) contrail ice particle size distributions (Schröder et al., 2000) suggest that Gamma functions can be used to represent the small, growth-dominated ice particle modes in contrail cirrus as well (see below).

The diffusional growth rate for a single ice particle reads $\dot{r} = \gamma/r$ (Pruppacher and Klett, 1997), with the growth rate factor $\gamma = \nu D n_{\text{sat}}(S - 1)$. Here ν is the volume of a water molecule in ice, $D(p, T)$ is the diffusion coefficient of water vapor molecules in air, and $n_{\text{sat}}(T)$ is their equilibrium number density at ice saturation (Marti and Mauersberger, 1993).

Two processes affect the size distribution in the absence of sedimentation. On the one hand, smaller particles grow

faster than larger particles according to the above single particle growth law, which tends to narrow the size distribution. On the other hand, turbulence tends to broaden the size distribution due to mixing of air parcels containing particles of different size. As ice particle modes with very narrow size distributions are not observed, it is plausible to assume that growth-induced narrowing is at least in part compensated by turbulent broadening. It would therefore be inaccurate to use the above growth law while neglecting the effect of turbulence.

We now derive an effective depositional growth law that empirically accounts for the effect of turbulent broadening. The normalized Gamma size distribution

$$f = \frac{\lambda^{\mu+1}}{\Gamma(\mu+1)} r^\mu \exp(-\lambda r), \quad \lambda = \frac{\mu+1}{\bar{r}}$$

is characterized by two parameters, the shape parameter, μ , and the scale parameter, λ . Since the latter is inversely proportional to the modal (or mean) radius of the particle population, \bar{r} , depositional growth that increases \bar{r} (decreases λ) increases the distribution variance. As guided by in situ observations of young contrails in which sedimentation was not the main factor affecting the ice particle size distribution, we keep μ constant and thereby include the counteracting effects of growth and turbulence.

We seek a growth law that (i) represents the correct depositional growth of the mean of the ice particle population and (ii) is consistent with the shape-preserving assumption $\mu = \text{const.}$ Requirement (i) implies that the mean number radius should evolve according to the physical diffusional growth law, i.e. $\dot{\bar{r}} = \gamma/\bar{r}$, yielding

$$\bar{r}(t) = \sqrt{\bar{r}_0^2 + 2\gamma t}, \quad (2a)$$

with \bar{r}_0 defining the initial mean radius. Requirement (ii) implies that individual crystals should grow such that f does not change its shape (App. A), i.e. the relative change in the size of each particle should be equal to that of the mean, $dr/r = d\bar{r}/\bar{r}$, or, using Eq. (2a),

$$\dot{r} = \tilde{\gamma} r, \quad \tilde{\gamma}(t) = \gamma/\bar{r}^2(t) = \gamma/(\bar{r}_0^2 + 2\gamma t), \quad (2b)$$

the solution of which reads

$$r(t) = \sqrt{r_0^2 + 2(r_0/\bar{r}_0)^2 \gamma t}. \quad (2c)$$

Equation (2c) shows that initially small ice crystals ($r_0 < \bar{r}_0$) grow at a smaller absolute rate than larger ones.

The terminal fall speed for spherical particles (Pruppacher and Klett, 1997) is given by

$$v_t(r) = \alpha r^2, \quad \alpha = 2\rho g/(9\eta), \quad (3)$$

where ρ is the mass density of bulk ice, g is the acceleration of gravity, and $\eta(T)$ is the air viscosity. The use of Eq. (3) is sufficiently accurate, given the approximate nature of our model.

The horizontal wind is related to the constant shear rate σ via

$$u(z) = \sigma(z - z_c), \quad (4)$$

where z_c marks the fixed altitude where u changes its sign. The constant vertical wind $w > 0$ points upwards while particles sediment downwards with $v_t > 0$.

After introducing the auxiliary 2-D radial ice crystal size distribution function $F(x, z, r, t) = r f(x, z, r, t)$, Eq. (1) is cast into the form

$$\partial_t F + \tilde{\gamma} r \partial_r F + u \partial_x F + (w - v_t) \partial_z F = 0. \quad (5)$$

Equation (5) is solved using the method of characteristics, stating that F is constant along the set of characteristic curves $\{x(t), z(t), r(t)\}$. The function $r(t)$ is given by Eq. (2c), $x(t)$ and $z(t)$ follow from integration of $\dot{x} = u$ and $\dot{z} = w - v_t$:

$$z(t) = z_0 + (w - \alpha r_0^2) t - \alpha (r_0 / \bar{r}_0)^2 \gamma t^2, \quad (6a)$$

$$x(t) = x_0 + (z_0 - z_c) \sigma t + (w - \alpha r_0^2) \sigma t^2 / 2 - \alpha (r_0 / \bar{r}_0)^2 \gamma \sigma t^3 / 3. \quad (6b)$$

The use of the actual (constant) supersaturation s in calculating the growth factor γ would dramatically overestimate the water vapor deposition rate on contrail ice crystals and hence the associated sedimentation flux. This is because in cloudy regions the ice crystals tend to drive the supersaturation to zero. We have performed numerous LES simulations of contrail development over four to five hours of integration time and compared the results with those from the analytical solution to correct this effect empirically. The correction consists of replacing the actual supersaturation by a tuned supersaturation, \tilde{s} , that is used to compute the growth factor $\gamma = \nu D n_{\text{sat}} \tilde{s}$. This approach yields realistic results at various contrails ages (Sect. 2.1.3). At any fixed T , \tilde{s} is approximately a linear function in s , $\tilde{s}(s, T) \approx \xi(T)s$, with $\xi \simeq 0.1$. The ξ -values decrease with increasing T ; for example, $\xi = 0.114$ at $T = 220$ K and $\xi = 0.088$ at $T = 225$ K.

We consider contrail cirrus at the beginning of the atmospheric dispersion phase and therefore choose an initial time $t_0 = 2$ min. (Hereinafter the simulation time, or contrail cirrus age, $t \geq 0$ counts from t_0 onwards.) The initial contrail area is centered at $x_c = 0$ and $z_c = 11$ km, corresponding to $p = 230$ hPa at midlatitudes. For given contrail thickness h_0 and width b_0 , the contrail covers the altitude range $z_c \pm h_0/2$ and the horizontal range $\pm b_0/2$, with $h_0 = 250$ m and $b_0 = 400$ m. Our choices for h_0 and b_0 are upper limit values that may be connected with large shear and high relative humidity, as indicated by Lidar measurements (Freudenthaler et al., 1995), but yield the desired average total initial ice particle number density $n_0 = \mathcal{N}/(h_0 b_0) = 30 \text{ cm}^{-3}$ and result in a fair agreement with LES results (Sect. 2.1.3). A

Table 1. Fixed parameters used in contrail cirrus simulations. The subscript 0 denotes an initial condition.

air pressure p	230 hPa
contrail center altitude z_c	11 km
time after contrail formation t_0	120 s
contrail thickness h_0	250 m
contrail width b_0	400 m
number of ice crystals \mathcal{N}	$3 \times 10^{12}/\text{m}$
dilution factor power law index δ	0.65
shape parameter μ	3
total simulation time	4 h
PDF data sampling time interval	10 min
number of horizontal and vertical grid points	80×40

constant updraft velocity, $w = 5$ cm/s, is prescribed counteracting particle sedimentation; it does not affect the microphysical evolution (Sect. 2.1).

The shape parameter of the Gamma function representing the unimodal contrail cirrus ice crystal size distributions (see above) is set to $\mu = 3$. This is consistent with in situ data (Schröder et al., 2000), as we have verified by fitting Gamma functions to the observed size distributions. Initially, we assume the same distribution f in each grid cell. The above initial conditions are summarized in Table 1 along with other parameters that are kept fixed during the simulations. Processing of ice particles in the wake vortex regime (before t_0) affects n_0 and the mean ice particle radius \bar{r}_0 . For this reason, we introduce variability in these initial conditions as well (Sect. 2.2.1).

The final solution of Eqs. (2c), (5), (6a) and (6b) reads in terms of the radial distribution moments $M^{(k)}$ ($k \geq 0$):

$$M^{(k)}(x, z, t) = \mathcal{D}(t) \int_{r_l(x, z, t)}^{r_u(x, z, t)} r^k f(r, t) dr, \quad (7a)$$

$$\mathcal{D}(t) = [t_0/(t + t_0)]^\delta,$$

with

$$f(r, t) = \frac{\lambda^{\mu+1}}{\Gamma(\mu+1)} r^\mu \exp(-\lambda r), \quad \lambda(t) = \frac{\mu+1}{\bar{r}(t)} \quad (7b)$$

in which the mode radius \bar{r} evolves according to Eq. (2a), as shown in App. A. The dilution factor, $\mathcal{D}(t)$, needs to be introduced in Eq. (7a), because we have neglected entrainment of cloud-free air by turbulent mixing in the analytical solution. In our model, this process merely acts to reduce the particle number concentrations and affects the maximum contrail cirrus width defined in Sect. 2.1.2.

In the simple case of plumes of passive tracers, turbulent mixing by constant eddy diffusivities would lead to a dilution parameter $\delta = 1$. In our case, δ may vary with the evolving contrail shape; for instance, a strongly sheared and vertically extended contrail offers a larger surface area toward the surrounding air leading to enhanced turbulent mixing. Here, we

used a fixed value $\delta = 0.65$. Other choices for δ could be accommodated by changes in the tuned supersaturation.

Ice particles sediment and are advected according to Eqs. (6a) and (6b). The radii of the subset of particles with $\{x_0, z_0\}$ at $t = t_0$ that are present at $\{x, z\}$ at $t > t_0$ are determined as follows. We eliminate z_0 from Eqs. (6a) and (6b) to yield a relationship $r_0 = r_0(x, z, t; x_0)$:

$$r_0^2(x, z, t; x_0) = \frac{1}{4q-1} \frac{6}{\alpha t} \left[\frac{x-x_0}{\sigma t} - (z-z_c) + \frac{wt}{2} \right] \quad (8a)$$

with $q = 1 + \gamma t / \bar{r}_0^2$. Reformulating Eq. (6a) yields a second relationship of the form $r_0 = r_0(z, t; z_0)$:

$$r_0^2(z, t; z_0) = \frac{1}{q\alpha t} [(z_0 - z) + wt]. \quad (8b)$$

The intersection between Eqs. (8a) and (8b) contains the corresponding r_0 -values. The actual radii $r(x, z, t)$ are then obtained from Eq. (2c), from which the lower and upper limit values r_l and r_u can be derived for use in Eq. (7a). The size distribution of ice particles at any given point for $t > t_0$ is therefore no longer represented by a full Gamma function due to the absence of particles with $r < r_l$ and $r > r_u$.

2.1.2 Microphysical and optical properties

The ice crystal number density, effective ice crystal radius, and ice water content follow from

$$n(x, z, t) = n_0 M^{(0)}, \quad (9a)$$

$$r_{\text{eff}}(x, z, t) = M^{(3)} / M^{(2)}, \quad (9b)$$

$$\text{IWC}(x, z, t) = 4\pi \rho n_0 M^{(3)} / 3, \quad (9c)$$

respectively, with the bulk ice mass density ρ . The ice water path is obtained by vertical integration:

$$\text{IWP}(x, t) = \int_{z_l(x, t)}^{z_u(x, t)} \text{IWC} dz. \quad (9d)$$

Eliminating r_0 from Eqs. (6a) and (6b) yields a relationship $z = z(x, t)$,

$$z(x, t; x_0, z_0) = z_0 + \frac{6q}{1+2q} \left[\frac{x-x_0}{\sigma t} - (z_0 - z_c) \right] - \frac{q-1}{1+2q} wt \quad (9e)$$

from which the cloud top and cloud base altitudes z_u and z_l can be calculated for use in Eq. (9d).

The shortwave optical extinction (essentially scattering) for spherical particles is calculated from

$$E(x, z, t) = \pi n_0 \mathcal{D}(t) \int_{r_l(x, z, t)}^{r_u(x, z, t)} Q_{\text{ext}} r^2 f(r, t) dr, \quad (10a)$$

with the Mie extinction efficiency, Q_{ext} , approximated by (van de Hulst, 1957)

$$Q_{\text{ext}}(r) = 2 - \frac{4}{\varrho} \left[\sin(\varrho) - \frac{1 - \cos(\varrho)}{\varrho} \right], \quad \varrho = 4\pi r(\kappa - 1) / \lambda_v. \quad (10b)$$

We choose $\kappa = 1.31$ for the real refractive index of ice and $\lambda_v = 0.55 \mu\text{m}$ for the wavelength of visible light. It is sufficient to apply Q_{ext} at one representative wavelength since for $r > 1 - 2 \mu\text{m}$, Q_{ext} rapidly approaches its limiting value of 2. The vertically integrated extinction defines the shortwave column optical depth τ :

$$\tau(x, t) = \int_{z_l(x, t)}^{z_u(x, t)} E dz. \quad (10c)$$

The maximum contrail cirrus width, $b_{\text{max}}(t)$, follows from

$$b_{\text{max}}(t) = [x_u(t) - x_l(t)] / \mathcal{D}(t), \quad (11)$$

with the upper (right) and lower (left) horizontal domain boundaries x_u and x_l determined by the characteristic equations. The effect of turbulent mixing is accounted for by the dilution factor. Equation (11) is a measure of the total contrail cirrus horizontal coverage and includes optically very thin portions, as entrainment of clear air will primarily create layers with very few (possibly large) ice particles (Heymsfield et al., 1998). Therefore, b_{max} will often be larger than the width actually visible by a ground-based observer or detectable by a satellite sensor. Furthermore, the true width will often be affected by the vertical extent of supersaturated layers, as addressed in Sect. 2.2.2. We note that in order to create probability distributions of contrail cirrus properties, every contrail was sampled homogeneously 80×40 times irrespective of its actual dimension.

Our solution is fully analytical and exactly conserves the domain-integrated ice crystal number density, corrected for dilution, at all times:

$$\mathcal{N} = \int_{x_l(t)}^{x_u(t)} \int_{z_l(x, t)}^{z_u(x, t)} \frac{n(x, z, t)}{\mathcal{D}(t)} dz dx = \text{const.} \quad (12)$$

All spatial integrations have been carried out using the trapezoidal rule. We have coded the above analytical solution into a standard Fortran subroutine, the contrail cirrus simulator CCSIM.

2.1.3 Model validation

We demonstrate the quality of our approximate solution by a comparison of CCSIM results to those obtained from an LES model that uses a two-moment bulk ice microphysical scheme (Spichtinger and Gierens, 2009). Figure 1 displays simulated distributions of ice water content after $t = 2 \text{ h}$. The LES run (left panel) was initialized with a realistic initial spatial distribution (300 m thickness, 150 m width) of contrail microphysical properties obtained from a simulation

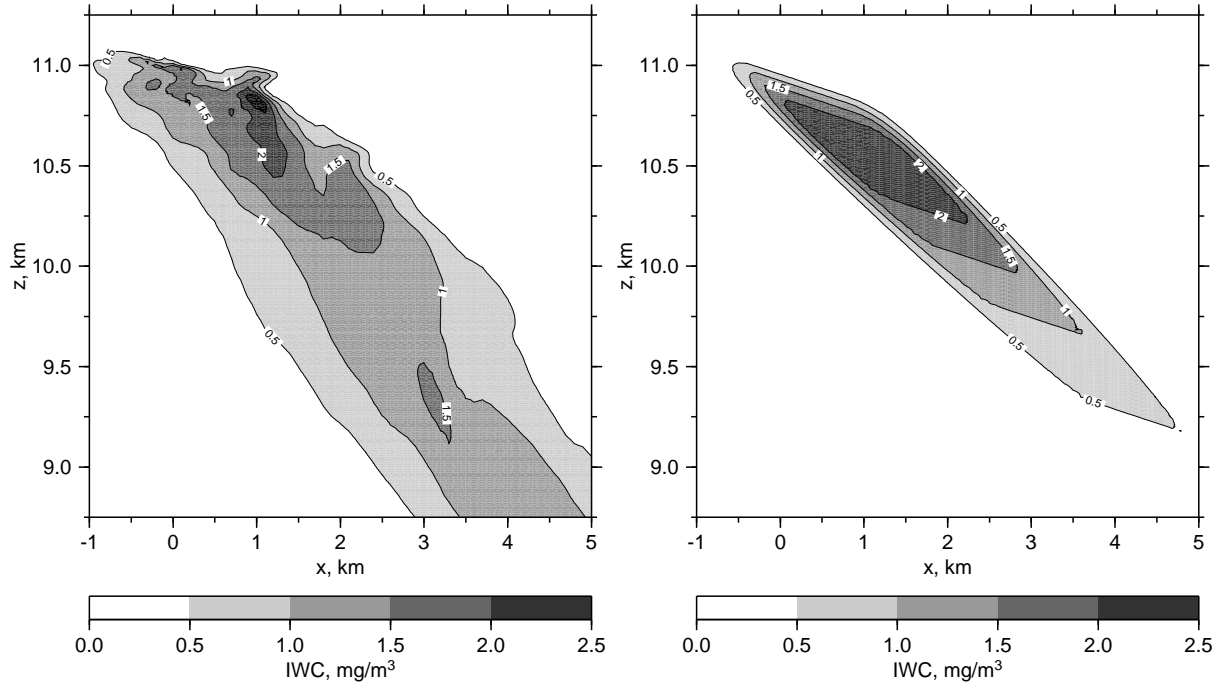


Fig. 1. Distribution of ice water content in an optically thin contrail cirrus cloud of age 2 h from (left) LES and (right) the CCSIM. In the LES, the initial IWC distribution was obtained from a simulation of the vortex phase evolution of the contrail and covered a region 300×150 m centered at $z = 11$ km and $x = 0$ km. In the CCSIM, the initial IWC was homogeneously distributed over a 200×100 m square area centered at the same location.

of the contrail development during the vortex phase (Unterstrasser et al., 2008). The ice particle size distribution f in the CCSIM (right panel) was initialized uniformly within a square area (200 m thickness, 100 m width). The LES uses a temperature lapse rate, while T is constant in the CCSIM. Otherwise, the initialization in both models is based on $\sigma = -0.001 \text{ s}^{-1}$, $w = 0$, $s = 0.15$ ($\tilde{s} = 0.017$), $T = 220$ K, and a total ice particle number density of $n_0 = 16 \text{ cm}^{-3}$. The resulting initial effective radius ($3 \mu\text{m}$), ice water content (0.9 mg/m^3), and extinction (0.5 km^{-1}) are consistent with aircraft data (Sect. 4.1). In the LES model, the cloud variables and T from above (as used in the CCSIM) were representative for the contrail core region, in which the air was almost saturated due to the previous ice particle growth.

Because the LES model captures turbulence and small-scale variability in supersaturation contrary to the CCSIM, the LES result in Fig. 1 shows more spatial structure. This has been verified by comparing LES results with and without homogeneous initial conditions (not shown); these two IWC fields look very much alike. Importantly, all essential features of the IWC distribution from the LES are reproduced by the CCSIM results. Among those are the magnitude of the peak IWC of $\sim 2 \text{ mg/m}^3$, the width, as well as the location and tilt of the region in which the IWC is largest. Despite the differences visible in the spatial distribution of IWC, the resulting ice water paths are similar.

Another difference is seen below ~ 9.5 km, where the LES contrail exhibits a fallstreak containing ice crystals with very large sizes (effective radii $> 80 \mu\text{m}$). This is likely caused by warmer temperatures at lower layers, since the LES prescribed a lapse rate of $\sim 8 \text{ K/km}$. Ice crystals falling into warmer supersaturated regions experience enhanced growth rates; as $v_t \propto r^2$, their sedimentation flux increases strongly with r . This feature is missing in the CCSIM, because T is assumed constant in the entire domain. Another cause for the generation of very large crystals in the LES might stem from the fact that the sedimentation of bulk mass and number concentrations are parameterized independently as a function of a mean ice particle radius. The respective sedimentation fluxes may become inaccurate in regions where both quantities are small.

The distribution of effective radii from the CCSIM is given in Fig. 2, showing the typical vertical layering of particle sizes caused by growth in a supersaturated and sheared environment, as also seen in LES simulations using size-resolved ice microphysics (Jensen et al., 1998, their Plate 3). Ice particle sizes from our LES model (not shown) are in good agreement with the CCSIM results in the high IWC-region.

The contour lines in Fig. 2 mark the contrail cirrus extinction. The region of maximum values $\sim 0.1 \text{ km}^{-1}$ is shifted to slightly higher altitudes relative to the high IWC-region and includes part of the cloud top where IWC is low. Small

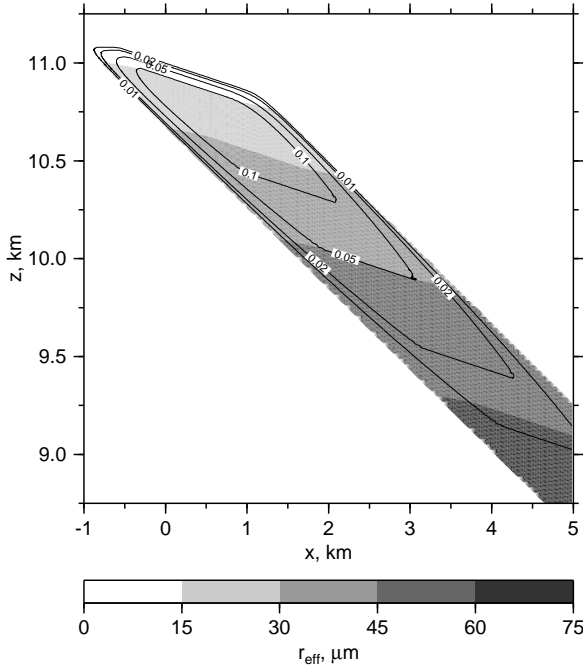


Fig. 2. Distribution of effective ice crystal radii and extinction at $0.55 \mu\text{m}$ wavelength (contour lines, in units of km^{-1}) from the CCSIM corresponding to the IWC distribution shown in Fig. 1 (right panel).

ice crystals with $r_{\text{eff}} \simeq 10 - 30 \mu\text{m}$ contribute most to the efficiency of the contrail cirrus to scatter visible light.

The probabilities of occurrence of corresponding cloud optical depths τ taken at different contrail cirrus ages are given in Fig. 3, both from (a) the LES and (b) the CCSIM. We use Eqs. (10a)–(10c) along with the exact ice crystal size distributions to compute τ in the spherical approximation in the CCSIM. In the LES, r_{eff} is parameterized from total crystal mass and number concentrations using assumed log-normal size distributions. From these values, τ is obtained from the parameterization of Ebert and Curry (1992), assuming ice crystals to be hexagonal columns with variable aspect ratio (spheres for maximum crystal dimension $< 7 \mu\text{m}$).

The general trend of the optical depth evolution is very similar in both cases. Generally, the initial LES distribution is characterized by more localized spots of high optical depth, contrary to the CCSIM which uses a uniform initial distribution. The influence of the initial conditions on the distribution of τ diminishes with cloud age. The mean τ decreases over time and optically thin cloud regions become more prevalent. The most marked difference between the two model distributions is seen at 30 min, with CCSIM showing a larger spread of high τ -values than the LES. This is caused by the differences in initialization (see above), as revealed by the solid curve with filled circles obtained from the LES using homogeneous initial fields, which is closer to

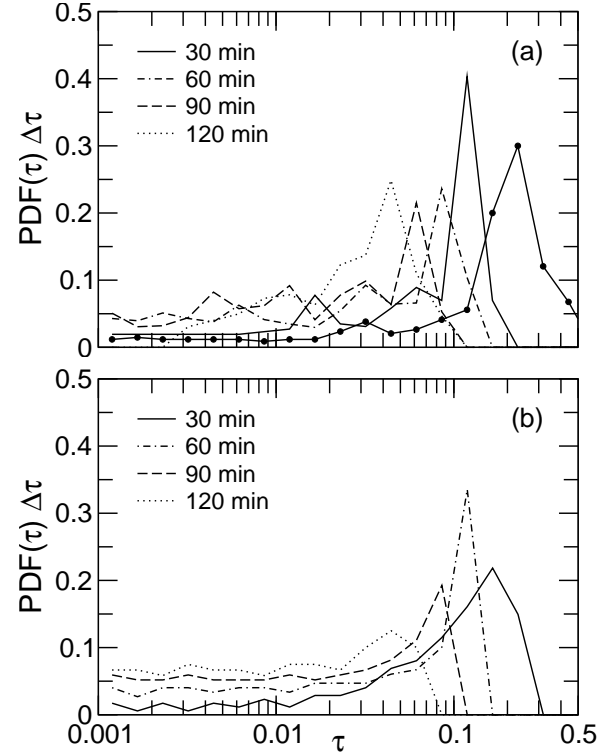


Fig. 3. Probability of occurrence of contrail cirrus optical depths from (a) LES and (b) the CCSIM at various cloud ages. Simulations as in Fig. 1, except the solid curve with filled circles in (a) that was obtained using a uniform initial distribution of cloud properties.

the solid curve from Fig. 3b. Differences in the mean τ at 30 min are due to initial differences $\tau_0 = 0.1$ (CCSIM) and $\tau_0 = 0.17$ (LES), mainly caused by the assumed initial contrail thickness of 300 m in the LES and 200 m in the CCSIM (chosen to be comparable to the vertical extent of the contrail core in the LES model). Given the differences in estimating τ in both models, the agreement between the major features of contrail cirrus evolution is judged as good.

Similar comparisons have been carried out over the entire range of temperatures, supersaturations considered in this work, and contrail ages up to 5 hours. We have generally found a fair agreement and regard the comparison shown in Figs. 1–3 as typical. Therefore, we believe that the CCSIM is an appropriate and sufficiently accurate tool to carry out statistical analyses of contrail cirrus properties.

2.2 Variability of cloud-controlling factors

2.2.1 Probability distribution functions

A probability distribution function $\text{PDF}(y)$ of the variable y is defined as the frequency of occurrence of y -data in the interval $[y, y + dy]$ normalized to unity. Our PDF describes

the atmospheric mean state and variability of y relevant to the evolution of contrail cirrus.

Below we define realistic PDFs of the key cloud-controlling factors T , s , σ , and the vertical extension d of ice-supersaturated layers (Fig. 4), as well as of the initial total number density n_0 and mean radius \bar{r}_0 of contrail ice crystals. Together these variables determine contrail cirrus properties. The employed PDFs are largely based on observations and help determine the most crucial factors impacting contrail cirrus optical depth and its variability. We will also employ uniform distributions of y around a mean value y_m (no variability), $\text{PDF}(y) = \delta(y - y_m)$.

Table 2. Parameters introduced to study the variability of cloud-controlling factors. LOW, BASE, HIGH are not uniquely defined but depend on actual combinations of these parameters.

case	T_m , K	σ_T , K	s_m	σ_m , s ⁻¹	d_m , km
LOW	218	1	0.05	0.002	0.5 ^a
BASE-C	214	2	0.15	0.004	1
BASE	218	2	0.15	0.004	1 ^b
BASE-W	222	2	0.15	0.004	1
HIGH	218	4	0.25	0.008	2 ^c
range	$[T_m \pm 3\sigma_T]$		[0,0.5]	[0,0.02]	[0.25,3] ^d

^a $d_{m1} = 100$ m, $d_{m2} = 750$ m.

^b $d_{m1} = 300$ m, $d_{m2} = 1150$ m.

^c $d_{m1} = 600$ m, $d_{m2} = 1750$ m.

^d Lower limit thickness given by initial contrail thickness h_0 .

The parameters defining a baseline as well as perturbed simulation cases are compiled in Table 2. We note that LOW, BASE, and HIGH do not represent unique cases; their actual definitions depend on whether uniform or non-uniform PDFs are used and whether one or several parameters from Table 2 are varied at a time. The parameter combinations associated with each case will be explained where appropriate.

Regional temperature variability can be approximated by Gaussian PDFs according to airborne measurements (e.g. Kärcher and Haag, 2004),

$$\text{PDF}(T) = \frac{1}{\sigma_T \sqrt{2\pi}} \exp \left[-\frac{1}{2} \left(\frac{T - T_m}{\sigma_T} \right)^2 \right]. \quad (13)$$

We fix the baseline mean temperature $T_m = 218$ K, which is a typical value at extratropical cruising altitudes near 220 hPa (Kärcher et al., 2009). For sensitivity studies, we assume T_m to be 4 K colder or warmer, corresponding to conditions at higher (200 hPa) and lower (260 hPa) flight levels. Our base case assumes a standard deviation of $\sigma_T = 2$ K. To capture the effect of enhanced or reduced variability in ice-supersaturated regions, e.g. dependent on gravity wave activity, σ_T will be varied by a factor of 2 in either direction.

Supersaturation with respect to the ice phase is observed at times in the upper troposphere (e.g. Jensen et al., 2001; Gettelman et al., 2006). We regard in situ measurements of rel-

ative humidity at the temperatures of interest as sufficiently accurate for our purpose. The corresponding PDFs are approximately exponential functions,

$$\text{PDF}(s) = \frac{1}{s_m} \exp \left(-\frac{s}{s_m} \right). \quad (14)$$

It is possible that the slope of the PDF in regions with higher T is steeper than in colder regions since superaturation relaxation time scales decrease with increasing T . Such a correlation effect is not captured by using the above PDF. We use the mean value of the PDF $s_m = 0.15$ as a value typical for regions in the flight levels where ice supersaturation occurs (Gierens et al., 1999). We vary s_m for sensitivity studies between the values 0.05, leading to a high frequency of occurrence of slightly supersaturated cases, and 0.25, where a small number of cases in the distribution tail already come close to the homogeneous freezing limit of cirrus formation.

Contrail spreading is mainly caused by vertical shear of the horizontal wind field. Spreading might also be induced by radiative heating and associated lateral expansion in a stably stratified environment, but given the rather high upper tropospheric wind shear, this contribution is likely to be of minor importance. A PDF of wind shear has been derived from European Centre for Medium Range Weather Forecasts (ECMWF) analysis data covering the North Atlantic flight corridor (Dürbeck and Gerz, 1996). We found that their PDFs are reproduced very well by a Weibull distribution,

$$\begin{aligned} \text{PDF}(\sigma) &= W(\sigma; c, \sigma_m) \\ &= \frac{c}{\sigma_m} \left(\frac{\sigma}{\sigma_m} \right)^{c-1} \exp \left[-\left(\frac{\sigma}{\sigma_m} \right)^c \right] \end{aligned} \quad (15)$$

with $c = 1.6$ and $\sigma_m = 0.003$ s⁻¹, corresponding to a change in u of 3 m/s per km of altitude change. We use a mean shear rate of $\sigma_m = 0.004$ s⁻¹ to define the base case. On the one hand, the ECMWF analysis may not capture cloud-scale shear rates, which are likely larger than assumed here. On the other hand, only the shear component perpendicular to the flight axis acts to spread the contrail cirrus width, so the mean shear values that actually cause the spreading may be smaller than those present at the scale of the cloud.

This view is supported by analyses of wind shear measured onboard the DLR research aircraft Falcon in the North Atlantic flight corridor (Unterstrasser, 2008). The corresponding cloud-scale PDFs are also fitted well by Eq. (15). These analyses yield $\sigma_m = 0.008$ s⁻¹; averaging the measured shear vector over flight directions resulted in $\sigma_m = 0.005$ s⁻¹. Furthermore, shear values of $0.002 - 0.004$ s⁻¹ are consistent with estimates inferred from contrail spreading rates measured using Lidar in conditions supporting contrail persistence (Freudenthaler et al., 1995). Therefore, we regard our choice of the baseline σ_m and its variations (low shear case with $\sigma_m = 0.002$ s⁻¹ and high shear case with $\sigma_m = 0.008$ s⁻¹) as appropriate.

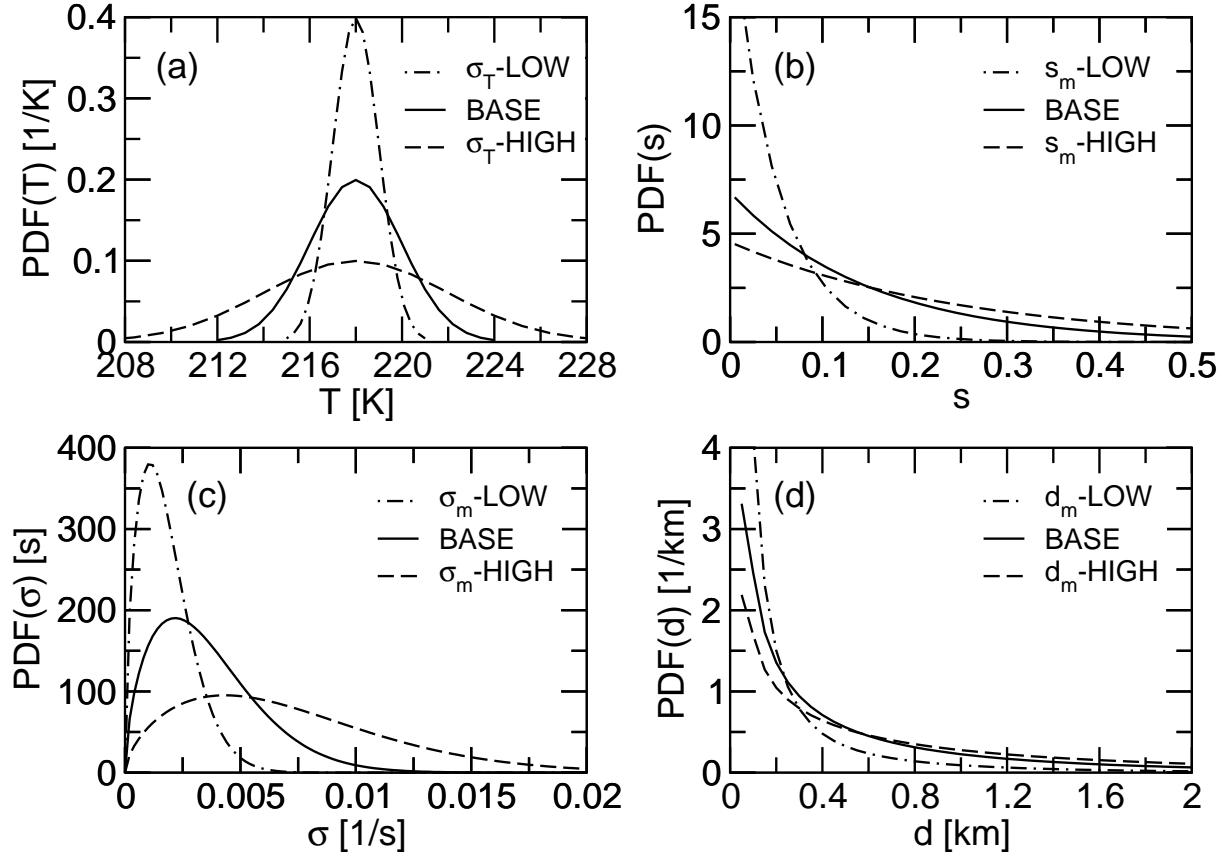


Fig. 4. Normalized probability distribution functions (PDFs) of the contrail cirrus-controlling factors (a) temperature variability, (b) ice supersaturation, (c) vertical shear of the horizontal wind, and (d) thickness of ice-supersaturated layers. The baseline (BASE) PDFs are given as solid curves, perturbed cases (LOW, HIGH) are indicated by dot-dashed and dashed curves, respectively. See Table 2 for details.

Statistics of the vertical thickness d of ice-supersaturated layers in the upper troposphere are available from radiosonde measurements. Measurements over northern Germany (Spichtinger et al., 2003) have been represented by a weighted sum of two Weibull distributions,

$$\text{PDF}(d) = c_0 W(d; c_1, d_{m1}) + (1 - c_0) W(d; c_2, d_{m2}), \quad (16)$$

with $c_0 = 0.75$, $c_1 = 0.7$, and $c_2 = 1$, yielding a mean thickness $d_m = 560$ m along with a large standard deviation of 610 m. Analyses of data taken in the Arctic (Treffeisen et al., 2007), over southern England (Rädel and Shine, 2007), and over Boulder, CO (Dowling and Radke, 1990) reveal larger annual mean values, $d_m = 1$ km (0.6–1.2 km), $d_m = 1.1$ km (0.9–1.3 km), and $d_m = 1.4$ km, respectively, with thicker (thinner) layers occurring in winter (summer). The latter two studies support the use of Eq. (16). One study reports a mean contrail formation layer thickness of 1.7 km (3.7 km) in the summer (winter) season over Fairbanks, Alaska (Stuefer et al., 2005). These latter values systematically exceed the above quoted supersaturated layer depths, as they include subsaturated conditions in which contrails form but are not persistent.

There are uncertainty issues with the radiosonde analyses. Only some of the above data sets have been corrected for dry biases. Nevertheless, even the corrected data may still have a dry bias. Furthermore, vertical resolution also affects d_m . In addition, variability in meteorological conditions may explain part of the differences in d_m -values inferred at different locations; this variability is also apparent in the data discussed by Kosarev and Mazin (1991), but no mean values were given. These data include locations in India, Indonesia, and Malaysia.

Regardless, the above d_m -values appear to be consistently smaller than average midlatitude cirrus cloud thicknesses of ~ 2 km derived from Lidar (e.g. Wang and Sassen, 2002) or satellite data (e.g. Lamquin et al., 2008). The radiosonde data indicate at times that two or more shallow supersaturated layers are vertically stacked, with subsaturated regions residing in between. Upward motion and resulting cooling of the air, which allows ice formation and supports cirrus persistence, would tend to increase the relative humidity also in the drier areas. This may explain cirrus thicknesses being larger than supersaturated layer depths prior to cirrus forma-

tion. Also, larger ice particles may fall substantial distances through subsaturated air before they completely sublimate. It is therefore plausible to vary the parameters of $\text{PDF}(d)$ such that the mean d_m falls within the noted limits of 0.5 km and 1.5 km, with $d_m = 1$ km representing our base case.

The number density n_0 of initial contrail ice crystals depends on a variety of aircraft and environmental parameters. To account for the variability of about a factor of 10 in n_0 caused by varying emission properties and plume dilution rates during contrail formation (Kärcher and Yu, 2009) and additional variations due to aircraft weight and engine configuration (Sussmann and Gierens, 2001), wake vortex dynamics (Lewellen and Lewellen, 2001), and ambient stability and turbulence (Unterstrasser et al., 2008), we use

$$n_0 = \frac{\mathcal{N}}{h_0 b_0} \left(\frac{1}{3} + 3 \times \text{RAN} \right) \varphi_n, \quad (17a)$$

where RAN is a random number uniformly distributed within the interval $[0, 1]$, $\mathcal{N}/(h_0 b_0) = 30 \text{ cm}^{-3}$ that follows from Table 1, and $\varphi_n \leq 1$ is a correction factor explained in the following.

Ice crystal concentrations can be dramatically reduced by adiabatic compression that results from the dynamically-induced downward motion of the vortex system. We have derived expressions for the surviving fraction of ice crystal number per unit flight path length, φ_n , as a function of T and s , employed here with slight modifications (App. B). By default, we set $\varphi_n = 1$ and examine the impact of corrections $\varphi_n(s, T) < 1$ in a sensitivity study (Sect. 3.3).

Finally, to account for the observed variability in initial mean ice particle size, we set

$$\bar{r}_0 [\mu\text{m}] = 0.5 + 1.5 \times \text{RAN}. \quad (17b)$$

Smaller n_0 does not necessarily result in smaller \bar{r}_0 , as the smallest particles sublimate faster than the larger ones, so the resulting size distribution mean radius might change very little or even increase (Kärcher and Burkhardt, 2008). Little is known about the evolution of the contrail ice particle size distribution in the vortex regime. Therefore, the currently available information on n_0 and \bar{r}_0 does not permit a more detailed description in terms of inhomogeneous PDFs.

The use of Eqs. (17a) and (17b) results in a marked variability in n_0 and \bar{r}_0 as suggested by in situ measurements. Schröder et al. (2000) report number concentrations in the range $100\text{--}300 \text{ cm}^{-3}$ taken around 10 min contrail age while Heymsfield et al. (1998) report values $10\text{--}100 \text{ cm}^{-3}$ in a contrail up to 1 h old. This trend in n_0 with contrail age is roughly consistent with the dilution law in Eq. (7a). As n_0 from Eq. (7a) reflects the latter data set of older contrails, we may underestimate the true mean n_0 suggested by the former data set.

2.2.2 Implementation details

To initialize the simulations, the spatial contrail cirrus domain has been discretized into 80 horizontal and 40 vertical

grid points. The domain size is time-dependent and equal to the region actually covered by contrail cirrus ice crystals (Sect. 2.1.2), assuming a typical tilted shape in the presence of shear (Fig. 2). The ice crystals change their location and size with time t according to Eqs. (2c), (6a), and (6b). The analytical solution is recorded at preselected contrail ages. At each preselected age, we divide the actual contrail cirrus domain again equidistantly into 80×40 grid points. This yields a sufficiently high spatial resolution, regardless of the cloud extension. Although in principle we could track particles of an arbitrary initial size at any location, it is sufficient to determine the smallest and largest particle radii that reach a given grid point to evaluate the local size distribution moments. The resulting physical variables are thus representative for the grid cell volumes surrounding each grid point.

Before any call to CCSIM, we check whether a contrail would form in the first place (Schumann, 1996), using an overall propulsion efficiency of 0.33. If the formation criterion is not met, the subsequent call is dismissed. However, our results are not sensitive to the fact that contrails do not form above $T \simeq 224 \text{ K}$, because such temperatures have low statistical weight in extratropical conditions.

The supersaturated layer thickness d is taken into account by limiting the analysis to a cloud base altitude of $z_l(x, t) = z_u(x, t) - d$ in order to calculate column-integrated or cloud-averaged contrail cirrus properties. This ignores the fact that sublimation of larger ice particles is not instantaneous, which may lead to a small underestimation of ice water path and column-optical depth. Limiting the contrail cirrus vertical extent to d has implications for the true contrail width, b , which we estimate from

$$b(t) = \frac{d}{z_u^{\max}(t) - z_l^{\min}(t)} b_{\max}(t), \quad d \leq z_u^{\max} - z_l^{\min} \quad (18)$$

because the cloud area is represented by a tilted column in the $\{x, z\}$ -domain (Fig. 1).

The lifetime of contrail cirrus depends on the synoptic situation (Kästner et al., 1999). Warming of air in regions of subsidence acts to dissipate clouds unless ice particles sediment into moister layers. Because synoptic conditions are variable, a wide range of contrail cirrus lifetimes exists, depending on time and geographical location (Burkhardt and Kärcher, 2009). Even in air supporting the persistence of ice clouds, the column optical depth of contrail cirrus will diminish after a certain time span, because either the particles sediment into dry layers or wind shear tears the cloud apart. Prolonged lifetimes of 6 hours (Heymsfield et al., 1998) or 17 hours (Minnis et al., 1998) may be caused by longwave heating (Jensen et al., 1998) or by secondary nucleation of ice crystals in conditions of sufficiently strong cooling.

We simulate contrail cirrus for up to 4 hours. This limited period is motivated by the fact that our simulations employ fixed atmospheric conditions and neglect ice nucleation, and should therefore be viewed as a compromise between possible shorter and longer lifetimes that occur in nature. We have

confirmed that our results do not change significantly when increasing the maximum contrail cirrus age by one hour, but the assumption of fixed atmospheric conditions becomes less reliable.

The construction of PDFs of contrail cirrus properties as defined in Sect. 2.1.2 requires sampling of a representative amount of data during each CCSIM simulation. We sample the CCSIM output every 10 min for 4 hours in the entire variable $\{x, z\}$ -domain. One complete simulation thus provides PDFs of a given set of contrail cirrus properties, each based on a maximum of $80 \times 40 \times (6 \times 4)$ data points for n , r_{eff} , IWC, and E ; $80 \times (6 \times 4)$ data points for IWP and τ ; and 6×4 data points for b .

We have made the assumption that turbulent mixing acts to dilute the contrail area homogeneously in the CCSIM (Sect. 2.1.1). Properties such as τ are correspondingly reduced in proportion to $\mathcal{D}(t)$. Our homogeneous sampling across the model domain irrespective of the actual contrail dimension gives each contrail a similar weight, independent of its coverage. If sampling according to coverage, young contrails would be sampled less.

The PDF of a contrail cirrus property y follows from

$$\text{PDF}(y) = \sum_{jT=1}^{NT} \sum_{js=1}^{Ns} \sum_{j\sigma=1}^{N\sigma} \sum_{jd=1}^{Nd} W_{jT,j\sigma,jd} \times \text{PDF}(y\{x, z, t\}), \quad (19a)$$

where $\text{PDF}(y\{x, z, t\})$ results from the temporal and spatial sampling as described above. Each sampling range, as given in Table 2, has been discretized into 20 equidistant bins. Each CCSIM simulation is driven by the input parameters T_{jT} (at a particular σ_T), s_{js} , $\sigma_{j\sigma}$, and d_{jd} that represent the center values of discretized T -, s -, σ -, and d -arrays, respectively, as well as by the random choices of n_0 and \bar{r}_0 . The bin center values are used to drive the simulations. Further, the probability (or relative weight)

$$W_{jT,j\sigma,jd} = \text{PDF}(T_{jT}) \text{PDF}(s_{js}) \text{PDF}(\sigma_{j\sigma}) \text{PDF}(d_{jd}) \times \Delta T \Delta s \Delta \sigma \Delta d, \quad (19b)$$

accounts for the contribution of each cloud-controlling factor according to its normalized PDF.

3 Basic studies

In this section, we examine the effect on $\text{PDF}(\tau)$ of different cloud-controlling factors acting in isolation (Sect. 3.1) or together (Sect. 3.2). The more significant a cloud-controlling factor, the larger is the resulting PDF variance. We include extreme scenarios in support of the generation very optically thin or thick contrail cirrus. We study the PDFs of ice water path, effective radius, width and other properties of contrail cirrus as well. A sensitivity study reveals the systematic changes of the optical depth distributions caused by aircraft-induced ice crystals losses in the wake vortex regime (Sect. 3.3).

3.1 Role of individual cloud-controlling factors

We first discuss the probability distributions (Fig. 5) of contrail cirrus optical depth, allowing for variability in only one of the cloud-controlling factors at a time to explore their individual role in determining $\text{PDF}(\tau)$. The other variables have been fixed (no variability) and set equal to their baseline mean value given in Table 2. All cases assume $T_m = 218$ K. For each cloud-controlling factor, 20 realizations have been computed using $\varphi_n = 1$ and random n_0 (average value 30 cm^{-3}) and \bar{r}_0 (average value $1 \mu\text{m}$). The resulting poor sampling creates some statistical aberrations in the PDFs which are, however, not important for our discussion.

Using $\sigma_T = 2$ K in case BASE (solid curve in Fig. 5a), we obtain the generic skewed distribution as already indicated in the single cloud case (Fig. 3). However, due to inclusion of the variability in n_0 and \bar{r}_0 , the peak of $\text{PDF}(\tau)$ has broadened and its variance has substantially increased. Changing the standard deviation from σ_T -LOW (dot-dashed) to BASE to σ_T -HIGH (dashed) tends to create more optically thick contrails.

The significant effect on $\text{PDF}(\tau)$ of changing variability in ice supersaturation s is shown in Fig. 5b. Allowing for enhanced probability of occurrence of high supersaturation values (s_m -HIGH) slightly increases the number of optically thick events. Enhancing the occurrence of low supersaturation values (s_m -LOW) leads to a marked lowering of the mean optical depth, because at times less water vapor is available for deposition on ice particles. Compared to Fig. 5a, the region $\tau = 0.001 - 0.1$ is more populated as a result of the exponential distribution of s giving low supersaturations a higher statistical weight.

Allowing variability in the vertical shear of the horizontal wind σ (Fig. 5c), the $\text{PDF}(\tau)$ in case BASE is quite similar to its counterpart in Fig. 5a. Altering this variability from σ_m -LOW to σ_m -HIGH leads to changes that are more significant than those caused by variation of σ_T . When enhancing (lowering) the frequency of occurrence of high shear values, the resulting contrail cirrus shifts to smaller (larger) optical depths. This is because in sheared conditions, each vertical cloud column contains fewer ice crystals.

In Fig. 5d, only the supersaturated layer thickness d is varied. Compared with Fig. 5a, the peak region is much narrower, resulting in the region $\tau = 0.001 - 0.1$ to be less populated. This indicates that variability in d has much less impact on $\text{PDF}(\tau)$ than variability in the other cloud-controlling factors. Furthermore, Fig. 5d shows that $\text{PDF}(\tau)$ is very insensitive to changes from d_m -LOW (0.5 km) to d_m -HIGH (2 km). This is because in many cases, the main contrail cirrus features evolve within ~ 1 km thick vertical regions and occasional fallstreaks only contribute to a minor degree to large τ -values. Hence, changes in the high d -tail of $\text{PDF}(d)$ do not matter much. Furthermore, much of the sensitivity against small layer depths is taken out since we imposed a minimum layer thickness of $h_0 = 250$ m. We add that fall-

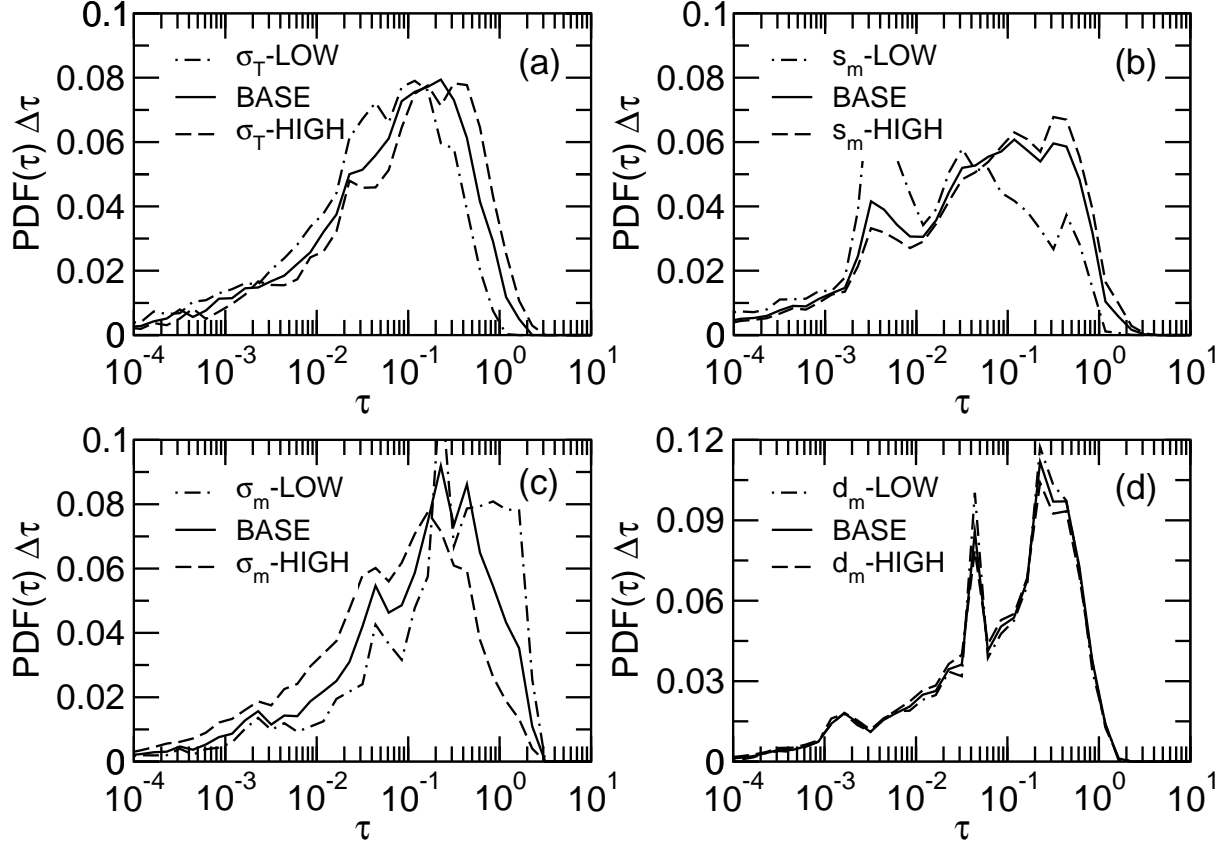


Fig. 5. Probability distributions of contrail cirrus optical depth obtained by imposing (a) variability in temperature T along with changes in temperature standard deviation σ_T (using constant baseline values for s_m , σ_m , d_m) and variability in (b) supersaturation s (constant σ_T , σ_m , d_m), (c) wind shear σ (constant σ_T , s_m , d_m), and (d) supersaturated layer thickness d (constant σ_T , s_m , σ_m) according to the cases shown in Fig. 4. All cases assume $T_m = 218$ K. Random sampling of initial number density n_0 (average value 30 cm^{-3}) and mean radius \bar{r}_0 (average value $1 \mu\text{m}$) is included with $\varphi_n = 1$. Spikes at intermediate optical depths are due to poor statistical sampling.

streaks in which most of the IWC is concentrated require at the same time large supersaturated layer depth, large supersaturation and weak wind shear; given the associated PDFs in Fig. 4, such cases are only seldomly realized.

3.2 Variability in contrail cirrus properties

The following studies assume BASE case variability in the variables T (via σ_T), d , and σ . In particular, we define two scenarios. First, the cases BASE, MIN, and MAX assume no variability in s . While BASE employs the baseline mean values s_m , σ_m , and d_m , we explore cases in which changes in s_m and σ_m support either low (s_m -LOW σ_m -HIGH, designated as MIN) or high (s_m -HIGH σ_m -LOW, designated as MAX) mean optical depth. Second, case FULL is identical to BASE but with exponential variability in s . This scenario is separated out because of the comparatively large influence exerted by variability in supersaturation (Fig. 5).

Probability distributions of τ , IWP, r_{eff} , and b are shown in Fig. 6. The cases MIN and MAX in Fig. 6a demonstrate that

contrail cirrus does not develop into optically thick cloud unless $s_m > 0.1 - 0.2$, i.e. in air masses undergoing persistent synoptic-scale cooling. The distribution mean optical depth becomes very sensitive to s_m below ~ 0.1 (MIN), decreasing very quickly with decreasing s_m relative to BASE, as we have studied by means of additional simulations (not shown). Optically thicker contrail cirrus evolves in regions with high s_m (MAX). High (low) σ_m -values aggravate this behavior in case MIN (MAX). However, the increase in the mean optical depth in MAX relative to BASE is limited by the fact that ice particles grow and fall faster out of the supersaturated layer. This may hint at the necessity of secondary nucleation occurring in long-lived contrail cirrus in order to maintain a high optical depth, a subject worthy of further studies. According to case MAX, the occurrence of contrail cirrus with $\tau > 3 - 5$ appears very unlikely. In nature, however, such cases may occur when contrails formed in cold conditions are advected into warmer supersaturated areas (possibly not allowing contrail formation).

The PDF of the ice water path (Fig. 6b) mirrors the behav-

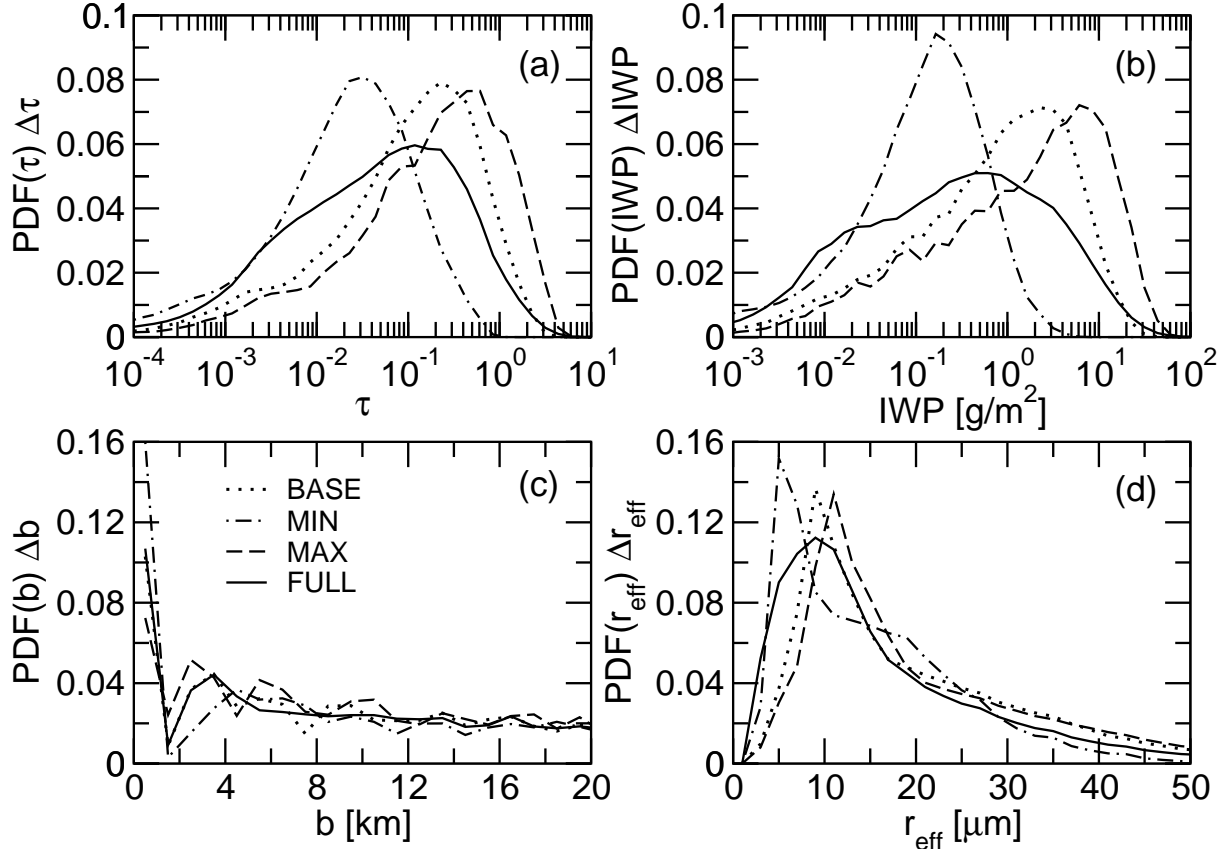


Fig. 6. Probability distributions of contrail cirrus (a) optical depth τ , (b) ice water path IWP, (c) width b , and (d) effective ice crystal radii r_{eff} . They have been obtained for $T_m = 218$ K by imposing simultaneous variability in temperature T , supersaturated layer thickness d , and wind shear σ according to their PDFs with $\sigma_T = 2$ K, $d_m = 1$ km, and $\sigma_m = 0.004 \text{ s}^{-1}$, and no variability in supersaturation s . Mean values s_m and σ_m vary according to Table 2. Random sampling of initial number density n_0 (average value 30 cm^{-3}) and mean radius \bar{r}_0 (average value $1 \mu\text{m}$) is included with $\varphi_n = 1$. Cases MIN (MAX) correspond to s_m -LOW σ_m -HIGH (s_m -HIGH σ_m -LOW). Case FULL is identical to BASE but additionally includes full variability in s according to its exponential PDF.

Table 3. Summary of contrail cirrus properties.

	τ		$E, \text{ km}^{-1}$		IWC, mg/m^3		IWP, g/m^2		$r_{\text{eff}}, \mu\text{m}$		$n, \text{ cm}^{-3}$		$b, \text{ km}$	
	mean	stdev	mean	stdev	mean	stdev	mean	stdev	mean	stdev	mean	stdev	mean	stdev
MIN	0.05	0.09	0.49	0.78	2.22	3.64	0.27	0.41	14.9	9.38	3.60	12.7	19.1	15.2
BASE	0.26	0.38	1.60	2.52	12.3	20.3	2.26	3.49	19.6	12.6	3.35	12.9	19.2	14.7
MAX	0.50	0.69	2.41	3.62	21.5	34.3	5.00	7.28	20.7	13.5	3.37	12.9	18.2	14.1
FULL	0.20	0.40	1.19	2.43	9.10	22.0	1.72	4.19	16.5	12.0	3.77	13.1	19.3	14.8

ior of $\text{PDF}(\tau)$, suggesting a limited variability in r_{eff} (see below). The PDFs of contrail cirrus width (Fig. 6c) are characterized by a large number of occurrences at widths < 1 km and a pronounced tail toward large b -values. (The sharp minima near $b = 1$ km, as well as the sharp peaks seen in some $\text{PDF}(r_{\text{eff}})$, are due to coarse temporal sampling because b , and r_{eff} , increase rapidly initially.) The fact that all PDFs look very much alike can be explained as follows. Shear induces spreading at a rate proportional to the contrail thick-

ness. The thickness grows with increasing supersaturation (or temperature) because of faster ice particle growth and sedimentation. Hence, the effect on b of lower shear values is partly compensated by an increased vertical extent. The overall effect is averaged out by variability in the supersaturated layer thickness, compare cases BASE and FULL.

The peak in $\text{PDF}(b)$ at small widths characterizes very young or only very weakly sheared narrow contrails that may not be detected in satellite imagery but are easy to probe

with aircraft. Many strongly sheared clouds with $b > 10$ km (60 – 65%) exist in all cases. The flat distribution tail may or may not contain detectable contrails. This tail may be associated with optically very thin contrails, as supported by additional simulations in which b was computed as the portion of contrail cirrus that exceeds a visibility or detectability threshold optical depth of 0.05. The resulting distributions (not shown) typically peak around $b = 4$ km and have much fewer wide contrail cirrus cases, the number of which decreases substantially from MIN to MAX.

The PDF of effective radius (Fig. 6d) peaks near $10 \mu\text{m}$ in case BASE, supporting the notion that contrail cirrus is on average composed of smaller particles than found in most natural cirrus clouds (Sect. 4.1). The peak r_{eff} -value increases (decreases) with increasing (decreasing) supersaturation, as expected. In all cases, a tail of r_{eff} -values reaching up to $\sim 50 \mu\text{m}$ is present, indicating that some regions contain few large ice crystals. In these regions, radiative forcing may differ significantly from that arising from the more localized core areas containing many small crystals.

The small r_{eff} -mode is clearly typical for young contrails, but may partly characterize older contrail cirrus as well. For instance, when contrails age their ice particles grow and fall to lower layers. If the supersaturated layer thickness is small, most of the contrail ice sublimates upon sedimentation but the small crystals with negligible fall speeds remain in the layer, presumably generating optically thin contrail cirrus. Whether these clouds become subvisible depends on the shear rate and on the ratio of timescales for depositional growth and sedimentation through the supersaturated layer.

Concerning case FULL, we confirm our earlier finding: the exponential form of $\text{PDF}(s)$ markedly enhances the number of optically thin contrails ($\tau < 0.1$) relative to BASE. A similar broadening towards small values is also seen in $\text{PDF}(\text{IWP})$, while the changes in $\text{PDF}(b)$ are small. Finally, variability in s enhances the probability to find r_{eff} in the range $4 - 8 \mu\text{m}$ at the expense of slightly larger values. We have repeated FULL without considering variability in d (not shown) and found $\text{PDF}(\tau)$ to be very similar, confirming the weak sensitivity to variations of d . We refer to Sect. 5.2 for more discussion on the impact of variability in d .

Table 4. Percentiles of contrail cirrus optical depth and fraction φ_v of subvisible contrail cirrus derived from the $\text{PDF}(\tau)$ shown in Fig. 6. The subvisible fraction is given by the cumulative PDF within $[0 < \tau \leq 0.02]$.

	10%	25%	50%	75%	90%	φ_v
MIN	0.001	0.005	0.02	0.06	0.12	0.51
BASE	0.003	0.023	0.10	0.31	0.61	0.23
MAX	0.006	0.043	0.20	0.60	1.18	0.17
FULL	0.002	0.007	0.04	0.16	0.44	0.39

Table 3 summarizes the mean values and corresponding

standard deviations of the properties discussed in Fig. 6 and adds similar information for contrail cirrus extinction, ice water content, and ice crystal number density. Table 3 highlights the fact that the PDFs of all contrail cirrus variables are characterized by standard deviations being comparable to or exceeding the mean values. This important finding has the following implications.

First, calculation of average optical depths is only possible if the full $\text{PDF}(\tau)$ is known. Neglecting optically very thin contrail cirrus as, e.g. in analyzing satellite measurements (Sect. 4.2), would yield artificially high mean values. Second, the description of processes that depend non-linearly on τ requires the use of the full variability in τ as described by its PDF. An example is the outgoing longwave radiation bias that is caused by neglecting cloud horizontal inhomogeneity, which is especially important for thin (e.g. $\tau < 1$) cirrus that occur in regions with temperatures much colder than the surface temperature (Fu et al., 2000).

As the PDFs are non-Gaussian, it is instructive to provide percentiles and median values in addition to mean and standard deviation, which we summarize in Table 4 for $\text{PDF}(\tau)$. In Table 4, we provide further information by the fraction of contrail cirrus that is invisible, $\varphi_v = \int_0^{\tau_v} \text{PDF}(\tau) d\tau$. The optical depth, τ_v , above which a cloud is visible depends on the photon wavelength, illumination conditions, viewing angle and distance to the observer, among other factors. We employ $\tau_v = 0.02$ as an approximate lower limit visibility threshold that is based on Lidar data and applies to ground-based observers (Sassen and Cho, 1992). According to Table 4, the subvisible contrail cirrus fraction is substantial and obviously inversely correlated with the mean optical depth. It rises from 17% to 51% from the most to the least favorable conditions in which contrail cirrus develops. Case FULL with $\varphi_v = 39\%$ is intermediate between MIN and BASE. Regions with very low τ include boundary areas and fall-streaks containing few particles, but also older contrail cirrus evolving in low supersaturation and/or high shear conditions. The possibility that contrails already start off with low τ is also included in the simulations due to the random choices of n_0 and \bar{r}_0 .

3.3 Sensitivity to wake processing of ice crystals

To examine the impact of ice crystal loss during the vortex phase on contrail cirrus optical depth (Sect. 2.2.1), we apply $\varphi_n < 1$ from App. B to estimate n_0 with the help of Eq. (17a). The factor $\varphi_n < 1$ reduces n_0 systematically with increasing T and decreasing s (Fig. B1). In this context, we report the range of initial ice water content and optical depth resulting from variability in n_0 and \bar{r}_0 . For example, in case BASE, setting $\varphi_n = 1$ leads to a range of initial values $\text{IWC}_0 = 0.01 - 5.7 \text{ mg/m}^3$ and $\tau_0 = 0.006 - 0.77$. Using $\varphi_n < 1$ yields $\text{IWC}_0 = 0.002 - 2.9 \text{ mg/m}^3$ and $\tau_0 = 0.001 - 0.4$, i.e. changes in the mean initial values by a factor of ~ 2 .

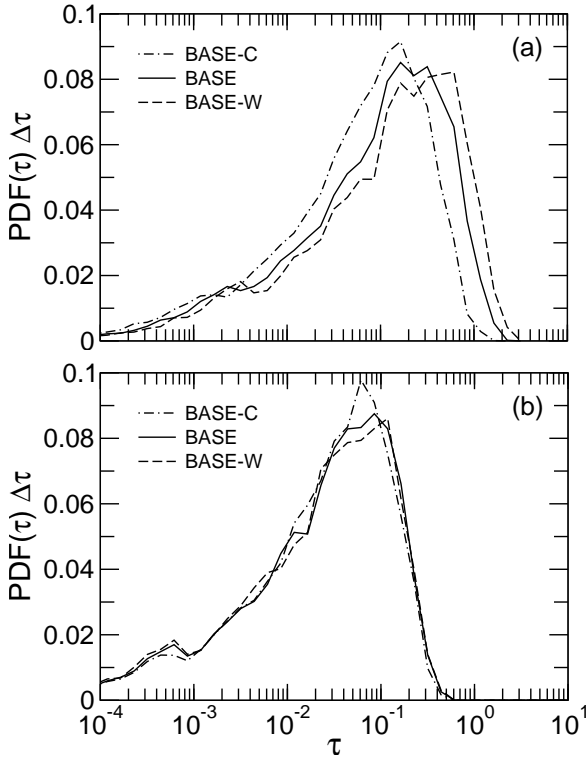


Fig. 7. Probability distributions of contrail cirrus optical depth. (a) Case BASE is identical to BASE from Fig. 6a but without variability in wind shear. The warm (BASE-W) and cold (BASE-C) cases assume mean temperatures $T_m = 222$ K and 214 K, respectively. (b) is identical to (a) but including ice crystal losses caused by adiabatic heating during the vortex phase, $\varphi_n < 1$.

The discussion below is not affected by variability in s and σ , for which reason we have neglected this variability in computing $\text{PDF}(\tau)$ shown in Fig. 7. Comparing the BASE cases from Fig. 6a and Fig. 7a reveals the impact of variability in wind shear on the PDF (which is included in Fig. 6a).

The effect on $\text{PDF}(\tau)$ of changing T_m from 214 K (BASE-C) to 218 K (BASE) to 222 K (BASE-W) at fixed $\sigma_T = 2$ K is shown in Fig. 7a. Warmer (colder) mean temperatures enhance (reduce) the mean optical depth, leaving the shape of $\text{PDF}(\tau)$ almost unchanged. We recall from Fig. 5a that increasing (decreasing) the variance of $\text{PDF}(T)$ has a very similar impact.

Figure 7b displays the same cases, but now evaluated with $\varphi_n(s_m = 0.15, T) < 1$. Considering the loss of ice crystals in this way approximately halves the mean optical depth and effectively removes the entire dependence of $\text{PDF}(\tau)$ on T_m . On the one hand, increasing (decreasing) T_m in case BASE-W (BASE-C) shifts the PDF to higher (lower) mean τ -values (Fig. 7a). On the other hand, φ_n strongly decreases (increases) when T_m increases (decreases), actually preventing optically thicker contrails from occurring in BASE-W

(creating more optically thicker contrails in BASE-C). These two systematic effects appear to cancel each other.

We expect that the systematic behavior of φ_n with s and T is likely to be realistically captured by the LES model in a qualitative sense. However, it remains unclear whether the resulting insensitivity of $\text{PDF}(\tau)$ on T_m , as suggested by Fig. 7b, is more realistic than the results from Fig. 7a. Effects caused by differences between larger and smaller sized aircraft have not been considered in the LES simulations determining φ_n . Furthermore, in the bulk microphysics approach employed in the LES, ice sublimation is parameterized using a prescribed log-normal particle size distribution that is not capable of developing the characteristic sublimation tail at low particle sizes (Kärcher and Burkhardt, 2008, their Appendix C). Hence, on the one hand, the function φ_n may not be accurate. On the other hand, our assumed initial concentrations n_0 in the range $10 - 100 \text{ cm}^{-3}$ from Eq. (17a) may underestimate the true mean n_0 , as noted in Sect. 2.2.1.

In sum, if φ_n -values significantly smaller than unity are common and therefore need to be introduced to initialize n_0 , it is likely that n_0 needs to be increased by a factor ~ 2 (still consistent with in situ data) in order to realize similar mean optical depths than discussed here, as these are in rough agreement with observations (Sect. 4.2). This issue warrants a detailed assessment. Ideal tools for such studies would be LES models that track the ice crystal size distribution, and a number of aircraft measurements.

4 Comparison with observations

Given the sparsity of in situ measurements of contrail microphysical properties (Sect. 1), a detailed validation of the results from Sect. 3 is not possible. However, it is fair to state that the simulated mean values of b , n , r_{eff} , IWC, E , and τ along with their variability ranges are within the confines of available experimental evidence, as detailed below.

4.1 Width and microphysical properties

Lidar measurements have been carried out in Southern Germany, suggesting widths between 1.5 km and 3 km for line-shaped persistent contrails (Freudenthaler et al., 1995) which were up to 30 min old. This range is at the lower boundary of the $\text{PDF}(b)$ shown in Fig. 6c, which is plausible as our simulations include contrails that spread over a much longer time. Tracking of contrails with Geostationary Operational Environmental Satellite (GOES) imagery with about 4 km resolution led to width estimates of 5 km and 10 km for contrails about 2 h and 4 h old, respectively (Duda et al., 2004). The cross-track width of a racetrack shaped contrail visible in GOES images was about 5 km (10 km) at 30 min (90 min) contrail age (Jensen et al., 1998). Apart from these data, width statistics of aging contrail cirrus based on observations are not available. We consider our results for contrail cir-

rus reasonable, as the model captures the main mechanisms affecting b . It is plausible that observed widths when estimated based on visibility are significantly smaller than the true widths from our model.

The mean ice crystal concentrations $n \approx 3.5 \text{ cm}^{-3}$ (Table 3) over the first 4 hours of age depend on the initial and mixing conditions (Sects. 2.1 and 2.2) and is in the upper range of ice crystal number densities observed in midlatitude cirrus mainly formed by homogeneous freezing (Kärcher and Ström, 2003; Gayet et al., 2004). More observations are required to better constrain this variable. The standard deviation $\sim 13 \text{ cm}^{-3}$ highlights the large in-cloud variability of n towards low concentrations caused by shear, sedimentation, and dilution. Variability in the temporal evolution of the dilution factor is not included in the model.

In our study, the mean ice crystal effective radii are $\sim 15 - 20 \mu\text{m}$ (Tab. 3). Those values are relatively small, because the available moisture must be shared among the large number of ice crystals (see above). In situ data typically show values $\sim 0.5 - 1 \mu\text{m}$ initially, increasing with contrail age to values $1 - 5 \mu\text{m}$ at up to $\sim 30 - 60 \text{ min}$ of age (e.g. Heymsfield et al., 1998; Poellot et al., 1999; Schröder et al., 2000; Febvre et al., 2009). This range of r_{eff} is roughly consistent with values derived from remote sensing and has been found to be much smaller than that from nearby cirrus (e.g. Betancor-Gothe and Grassl, 1993; Duda and Spinhirne, 1996). Two such studies tracked contrail clusters for a longer time (up to $\sim 5 \text{ hours}$) and inferred $r_{\text{eff}} \simeq 10 - 30 \mu\text{m}$ (Duda et al., 2001) and $r_{\text{eff}} \simeq 10 - 15 \mu\text{m}$ (Duda et al., 2004), roughly consistent with CCSIM. This is still below effective radii in midlatitude semi-transparent cirrus of $15 - 30 \mu\text{m}$ inferred from TIROS-N Observational Vertical Sounder (TOVS) satellite data (Stubenrauch et al., 2004) or lower than the range of 25 – 75 percentiles $13 - 35 \mu\text{m}$ (Gayet et al., 2004) and the range of mean values $20 - 40 \mu\text{m}$ (Liou et al., 2008) from various in situ measurements. It must be kept in mind that differences in retrieval methods, measurement techniques, and ice crystal habits render an exact comparison of r_{eff} difficult.

Small r_{eff} -values may at times be associated with subvisual contrail cirrus, in particular in conditions with small supersaturation, high wind shear, and small supersaturated layer depths (Sect. 3). This supports the hypothesis that the existence of nearly invisible ($\tau < 0.05$) ice crystal layers at the tropopause could be due partly to remnant contrail particles (Smith et al., 1998). The hypothesis was based on radiance observations during the Subsonic Aircraft Contrail and Cloud Effects Special Study (SUCCESS). The estimated parameters ($r_{\text{eff}} \simeq 7 \mu\text{m}$, IWP $\simeq 0.15 \text{ g/m}^2$) are close to those from our model case MIN, where the PDF(r_{eff}) peaks at $5 \mu\text{m}$ (Fig. 6a, Table 3).

Ice water content in young contrails has occasionally been measured onboard aircraft, indicating a range $2 - 6 \text{ mg/m}^3$ at temperatures $\sim 218 \text{ K}$ (Spinhirne et al., 1998; Schröder et al., 2000; Febvre et al., 2009). One study reports values of up to

18 mg/m^3 at 236 K (Gayet et al., 1996). An empirical relationship for IWC in persistent contrails and cirrus based on in situ measurements (Schumann, 2002) and a climatology of average IWC for thin cirrus derived from Lidar and Radar data (Wang and Sassen, 2002) both yield $\sim 3 - 4 \text{ mg/m}^3$ at 218 K . Gayet et al. (2004) report $3 - 18 \text{ mg/m}^3$ as a range of 25 – 75 percentiles (median value 8 mg/m^3) within typically $220 - 230 \text{ K}$. The small amount of available IWC data for contrail cirrus, the probing of small and spatially heterogeneous areas in contrails at high aircraft speeds ($> 150 \text{ m/s}$), and shortcomings of some cloud probes preclude final statements on the representativity of these data. Referring to Table 3, the simulated mean IWC values for contrail cirrus are $\sim 2 - 20 \text{ mg/m}^3$ at $218 \pm 2 \text{ K}$ bracketing the least to the most favorable conditions; in case BASE the mean is $\sim 2 - 3$ times higher than suggested by the observations. This may be caused by the relatively young age of observed contrails relative to the simulated contrail cirrus. The latter simply have had more time to deposit a greater amount of water vapor from supersaturated surroundings. The large simulated standard deviations of IWC are consistent with the large variability in IWC generally observed in cirrus, regardless of temperature (Schiller et al., 2008).

Few direct measurements of extinction at visible wavelengths are available. Sussmann (1999) reported an extinction profile of a young contrail in the early vortex regime (plume age 20 s) derived from an aerosol-lidar backscatter profile, with peak values of 3 km^{-1} (70 km^{-1}) for ice particles in the left (right) vortex. Febvre et al. (2009) report the optical extinction of young contrails (age $< 20 \text{ min}$) to be $0.29 - 0.48 \text{ km}^{-1}$ and that of a nearby cirrus cloud 0.3 km^{-1} based on aircraft measurements. Atlas et al. (2006) present Lidar measurements of contrails developing massive fall-streaks and report values in the range $0.5 - 2 \text{ km}^{-1}$. Typical extinction values for midlatitude cirrus appear to be in the range $0.1 - 0.5 \text{ km}^{-1}$ based on more extensive Raman Lidar measurements (Ansmann, 2002), or $0.2 - 1.2 \text{ km}^{-1}$ based on aircraft measurements (Gayet et al., 2004). Our calculated mean values for contrail cirrus range between $0.5 - 2.4 \text{ km}^{-1}$ (Table 3), exceeding those of measured cirrus because contrail cirrus contain more and smaller ice crystals consistent with the comparison of simulated r_{eff} and n with cirrus data.

These comparisons reveal substantial differences between the properties of contrail cirrus and extratropical natural cirrus up to at least 4 hours of age. It is not known in which circumstances (if at all) microphysical properties and the radiative response of contrail cirrus and natural cirrus converge. The properties of the simulated older contrail cirrus lie between those of line-shaped contrails and natural cirrus. Therefore, the CCSIM results are consistent with the observations.

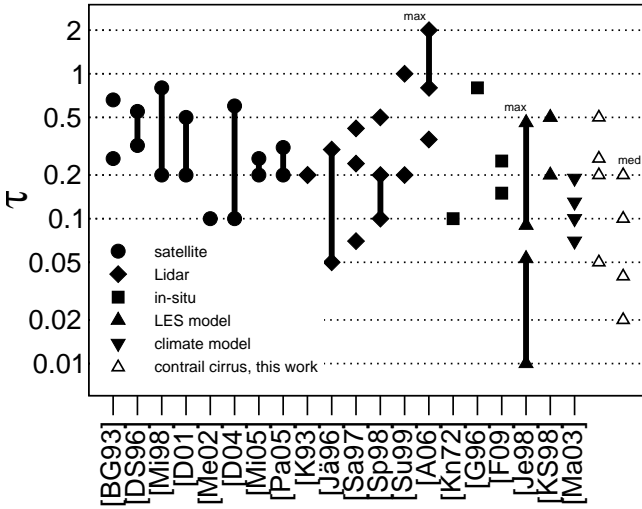


Fig. 8. Compilation of visible optical depths for observed and modeled persistent, line-shaped contrails from various sources. Data points represent mean values except when labelled “max” (maximum) or “med” (median). Open symbols are values for contrail cirrus from Tables 3 and 4. Several issues preclude a full intercomparability of the data, see text.

4.2 Optical depth

Figure 8 presents a compilation of visible optical depth data (symbols) for persistent, line-shaped contrails from various sources. These sources include satellite observations (Betancor-Gothe and Grassl, 1993; Duda and Spinhirne, 1996; Minnis et al., 1998; Spinhirne et al., 1998; Duda et al., 2001; Meyer et al., 2002; Duda et al., 2004; Minnis et al., 2005a; Palikonda et al., 2005), Lidar measurements (Kästner et al., 1993; Jäger et al., 1996; Sassen, 1997; Sussmann, 1999; Atlas et al., 2006), and aircraft in situ measurements (Knollenberg, 1972; Gayet et al., 1996; Febvre et al., 2009); models comprise LES (Jensen et al., 1998; Khvorostyanov and Sassen, 1998) and global climate model studies (Marquart, 2003; Ponater et al., 2002; Marquart et al., 2003), as well as the present approach (open symbols). Ranges of data from longer-term studies, time series, or sensitivity studies are given as symbols connected with lines.

Most of these data cannot directly be compared with each other for a number of reasons. Lidar optical depths are given at a particular photon wavelength, while many satellite data or models report values that are representative for a broader wavelength band in the visible region. This generates a scatter in optical depth values we have not accounted for. In situ and most Lidar studies represent single values or averages over single observations. Space-borne sensors are not capable of detecting very thin contrails (Sect. 5.1). Although possible, it is unlikely that contrails with $\tau < 0.01 - 0.1$ have frequently been probed by Lidar or aircraft, as they are hard to track visually, and such measurements are usually carried

out in conditions supporting strong contrail growth. Therefore, most observations are expected to be biased towards high optical depths.

Some satellite data cover extended areas and a full annual cycle, and therefore probe a variety of environmental states, while most other measurements and model results represent specific meteorological conditions or a narrow range of contrail age. Sussmann (1999) describes Lidar data taken in a very young contrail in the vortex regime (age 6 – 77 s), with the high value representative for contrail ice in the primary wake (wingtip vortices) and the low value characterizing ice in the secondary wake, which was separated from the former by a vertical gap roughly 50 – 100 m deep. This illustrates the high variability in initial contrail properties. Simulated contrails include cases in a strongly sheared environment (Jensen et al., 1998) that resulted in relatively low average optical depth, consistent with our case MIN. Optical depths reported by Betancor-Gothe and Grassl (1993) assume IWC of 4 (10) mg/m³ for the lower (upper) data point in Fig. 8. Kästner et al. (1993) discuss local enhancements above an aerosol/cirrus background signal. We have converted the infrared optical depths at a wavelength of 10.95 μ m reported by Duda and Spinhirne (1996) to visible optical depth by multiplying with a factor of 2, which is a reasonable approximation (Fu and Liou, 1993).

The climate model results for 1992 air traffic stand out, because they represent the only climatological annual mean values currently available for several regions (Marquart, 2003). By design, only line-shaped contrail optical depths have been estimated. Mean values in descending order are averages over Thailand, Japan, the U.S.A., and Western Europe. They show considerable regional variability and cover tropical regions, whereas all other references except Duda et al. (2001) exclusively report midlatitude measurements. The corresponding global mean figure ~ 0.1 formed the basis of previous global radiative forcing estimates for line-shaped persistent contrails (Sausen et al., 2005; Forster et al., 2007; Lee et al., 2009).

The climate model mean optical depths 0.07 – 0.19 from Fig. 8 have been obtained by averaging the simulated PDFs above the assumed visibility threshold $\tau_v = 0.02$; these PDFs include few peak values up to ~ 3 and many optically very thin contrails (Marquart et al., 2003). Hence, mean values would be lower had they been calculated from the full distributions. This may suggest a low bias. It is possible that the climate model mean values lie below the majority of observed mean values simply because the latter have a high bias. Given the above intercomparison issues, a more detailed analysis of climate model results is left for future work.

The open symbols in Fig. 8 are the CCSIM results obtained in the present study, including many contrail cirrus that are presumably not included in the other observational data sets. We show the mean and median optical depths for the cases given in Tables 3 and 4, recalling the large vari-

ances of the associated $\text{PDF}(\tau)$. The CCSIM values have been obtained only for a narrow, yet typical, temperature regime around 218 K, while other measurements and the climate model (see below) consider a much wider range of contrail occurrence, typically 208 – 228 K. The CCSIM results include the highest observed optical depth values in the distribution tails. The CCSIM mean values in case MIN are smaller than those derived from many of the measurements (0.1 – 0.5), which is to be expected due to the high bias of the latter.

5 Case study over the continental U.S.A.

Statistics of observed line-shaped persistent contrail optical depth covering extended areas and time spans of several months are only available from satellite observations, which is their major advantage over other measurements. We make a first attempt to compare $\text{PDF}(\tau)$ from our model with that of satellite data (Sect. 5.1). Additional contrail cirrus properties obtained from the model that could not be measured in the case study are discussed (Sect. 5.2).

Line-shaped persistent contrail optical depth has been derived over the continental U.S.A. (Palikonda et al., 2005). They applied an automated image processing algorithm (Mannstein et al., 1999) to radiances measured by the Advanced Very High Resolution Radiometer (AVHRR) on two different satellites during the period January–December 2001. The analysis method exploits the linear shapes of most contrails and large (compared to most natural cirrus) emissivity differences between two channels in the infrared. The analysis results were limited to satellite viewing angles $< 50^\circ$ and contrail optical depths $\tau \leq 1$. To deal with the detection issues (Sect. 4.2), a subjective error analysis was carried out. Correction factors were derived based on the detailed inspection of the automated algorithm's performance using 3 randomly selected days during 3 different months. Palikonda et al. (2005) derived optical depth statistics for each of the twelve months based on the corrected and diurnally adjusted detection frequencies for both satellites. The monthly mean probability of occurrence of τ was found to be remarkably consistent (within $\sim 20\%$), as were the statistics from both sensors.

For use in the CCSIM, meteorological conditions were taken from the National Oceanic and Atmospheric Administration (NOAA) Rapid Update Cycle (RUC) numerical weather analyses to constrain the CCSIM. The RUC assimilates a variety of input streams to produce hourly analyses and forecasts at relatively high spatial resolution over the continental U.S.A. (Benjamin et al., 2004). RUC temperature and wind data at the 250 hPa level were used to estimate the annual averages and variability of temperature and wind.

The results are $T_m = 223$ K and $\sigma_T = 4$ K for the Gaussian $\text{PDF}(T)$, and $\sigma = 0.0025 \text{ s}^{-1}$ and $c = 1.2$ for the Weibull-type $\text{PDF}(\sigma)$. In the context of the basic studies from

Sect. 3, these values represent rather warm conditions with low average wind shear, but large variability (broad PDFs) in both quantities. As for the relative humidity, the 2001 RUC data are not realistic (Minnis et al., 2005b). To have a more reasonable estimate of supersaturation distributions, we have used the exponential $\text{PDF}(s)$ and varied the mean value to obtain a best estimate of $\text{PDF}(\tau)$ that compares well with the observed PDF, yielding $s_m = 0.1$. The double-Weibull statistic with a mean $d_m = 1$ km has been used for the thickness of supersaturated layers. Finally, we have restricted the sampling of CCSIM data to $\tau \leq 1$ similar to the measurements and have used the same binning on a linear grid as Palikonda et al. (2005) to distribute the optical depths. The overall properties of the PDFs are comparable to case FULL (Fig. 6 and Tables 3 and 4), because of the large variability in cloud-controlling factors needed to represent the variety of conditions of an entire year of observations over the U.S.A.

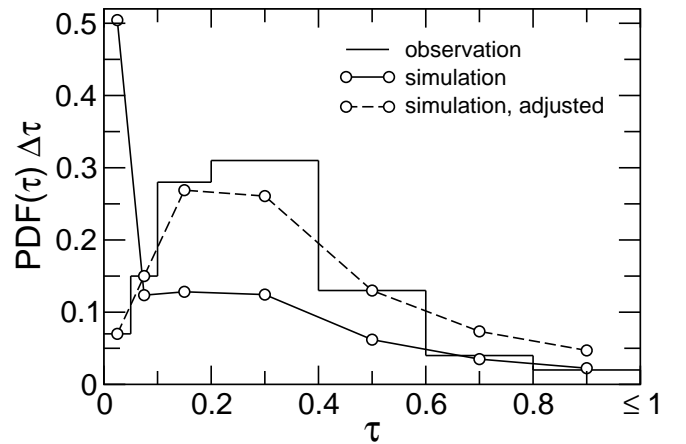


Fig. 9. Observed PDF of line-shaped persistent contrails (stepped lines), simulated PDF (solid curve) and simulated PDF adjusted with empirical detection efficiencies in the first two optical depth bins (dashed curve). Open circles mark the bin center values. For details, see text.

5.1 Statistics of optical depth

We have averaged the monthly data from Palikonda et al. (2005) to obtain a normalized $\text{PDF}(\tau)$ representative of 2001 (stepped lines in Fig. 9) that serves as a basis for the comparison with the CCSIM result (solid curve). In the model, every portion of contrail cirrus is sampled regardless of its optical depth. Compared to the observed PDF, a much larger number of occurrences is predicted in the first bin with $\tau \leq 0.05$. The shape of the simulated PDF agrees quite well with the observations in the region $0.6 < \tau \leq 1$, but underestimates the occurrences significantly within $0.1 < \tau \leq 0.6$. A small number of contrail cirrus with $\tau > 1$ (comparable to those present within $0.6 < \tau \leq 0.8$) are also predicted but are not included.

We recall a few limitations of contrail detection from passive space-borne remote sensors. There are no simple criteria describing the detection efficiency (fraction of cloud events detected) of contrails; for instance, linear contrails at least 10 km long with $b = 1 - 2$ km and $\tau > 0.5$ are easily observable, but may also be confused with natural cirrus of similar shapes resulting in false detections (Mannstein et al., 1999; Meyer, 2000; Minnis et al., 2005a). The detectability is poorer and generally not as well quantified for narrower structures, spreaded contrail cirrus, smaller optical depths, or contrail cirrus that has spread to more than few km wide. Furthermore, the quantification of contrail occurrence depends on factors such as concomitant cloudiness, sensor viewing angle, misidentification with natural cirrus, and ultimately on human perception. Those factors manifest themselves in tunable false alarm rates and introduce a subjective component into data analyses, which is difficult to consider in any comparison exercise.

Optically very thin contrails cannot be observed using passive space-borne remote sensing, affecting many detection schemes for cirrus identification (e.g. Wylie et al., 1995; Roskovensky and Liou, 2003; Krebs et al., 2007; Ackerman et al., 2008). References for the threshold optical depth above which a noticeable fraction of cloud is identified typically indicate values in the range 0.1 – 0.4. Efficiencies of detection decrease with decreasing τ because of smaller contrast between the contrail and background radiances and depend on the cloud physical properties and local situation during the measurement. Therefore, even for a specific sensor, a unique estimate of detection efficiencies cannot be given. The problem becomes even less well-defined when attempting to ascribe detection efficiencies to annual average data such as those from Fig. 9.

Regardless, for a meaningful comparison of simulated and observed PDF(τ), a correction accounting for the undersampling of optically thin contrails needs to be introduced. We proceed pragmatically by deducing empirical detection efficiencies ϵ as follows. In a first step, the observed and simulated data points in τ -bin 1 ($0 < \tau_1 \leq 0.05$) are used to derive $\epsilon(1) = P_{\text{obs}}(\tau_1)/P_{\text{sim}}^{(0)}(\tau_1) < 1$, the superscript denoting the original PDF. Next, the simulated data point in bin 1 is corrected to the observation, $P_{\text{sim}}^{(1)}(\tau_1) \equiv P_{\text{obs}}(\tau_1)$. The resulting updated PDF is re-normalized to unity, which increases the probabilities in all other bins accordingly; this results in a first estimate, $P_{\text{sim}}^{(1)}(\tau)$. Finally, it is checked whether $\epsilon(2) = P_{\text{obs}}(\tau_2)/P_{\text{sim}}^{(1)}(\tau_2) < 1$. If true, the procedure is repeated yielding a new normalized estimate $P_{\text{sim}}^{(2)}(\tau)$, etc. When $\epsilon(i) \geq 1$, the final adjusted PDF is given by $P_{\text{sim}}^{(i-1)}(\tau)$.

Detection efficiencies inferred in this way are termed empirical since they incorporate any difference in contrail properties between the observations (detecting line-shaped objects) and simulations (simulating very narrow and very broad contrail cirrus that remain undetected). Further, the efficiencies depend on differences in the way the PDF is sam-

pled in the model and in the observations.

The result of the above analysis is given as a dashed curve in Fig. 9. Adjustments have been made for the first two bins, yielding $\epsilon(1) = 0.14$, $\epsilon(2) = 0.65$, and $\epsilon(i) = 1$ for $i = 3, \dots, 7$. Multiplying the simulated PDF (solid curve) with these efficiencies and summing up the probabilities leads to the important finding that the satellite measurements have missed 47% of line-shaped persistent contrails of all optical depths.

These efficiencies relate to this specific observational case. Changes in the cloud-controlling factors, mean values and variance of the T -, σ -, s -, and d -distributions, led only to more optically thin contrail cirrus, worsening the agreement in the range $0.2 < \tau \leq 0.4$. The largest uncertainty is related to supersaturation (Sect. 3). We have found that the use of a uniform PDF(s) with $s_m = 0.08$ yields almost the same adjusted PDF(τ) as shown in Fig. 9, but with $\epsilon(1) = 0.20$ and $\epsilon(2) = 0.74$. This rather drastic change in variability of what we have identified as the most important cloud-controlling factor causes only little change in ϵ , because the respective τ -values are merged into relatively broad bins (compared to the progressively increasing logarithmic bin sizes). The remaining differences between the observed PDF and the adjusted CCSIM results may point to problems of the model to predict contrail cirrus at $\tau > 0.2$ or indicates unspecified uncertainty in the observation caused by the subjective error analysis (see above).

A large fraction of contrail cirrus is optically thin but those contrails are usually neglected in estimating the climate impact. The latter not only depends on the optical depth but also on the lifetime of contrails. Optically thin contrails developing in regions with low supersaturation may be short-lived and contribute little to the overall radiative impact. However, lifetimes of contrails developing in strongly supersaturated areas may be optically thick and also short-lived due to more rapid sedimentation. Furthermore, long-lived contrail cirrus may be optically thin due to high wind shear or low temperatures. Therefore, neglecting optically thin and subvisible contrails may lead to an underestimation of the total radiative impact.

5.2 Statistics of width and microphysical properties

We provide in Fig. 10 best estimate model PDFs (solid curves) for contrail cirrus parameters IWP, IWC, E , b , and r_{eff} over the continental U.S.A. that have not been measured by Palikonda et al. (2005) but which are consistent with the observed PDF(τ) and the RUC analyses. The PDF(τ) repeats the one from Fig. 9 (solid curve) on a logarithmic scale and additionally includes all events with $\tau \geq 1$. The PDF of IWP is significantly broader than that of IWC and differs from that of τ . This indicates that the optical depth distribution cannot simply be estimated from a PDF of IWC assuming a constant contrail cirrus thickness and effective radius. Instead, the vertical distributions of these cloud variables are

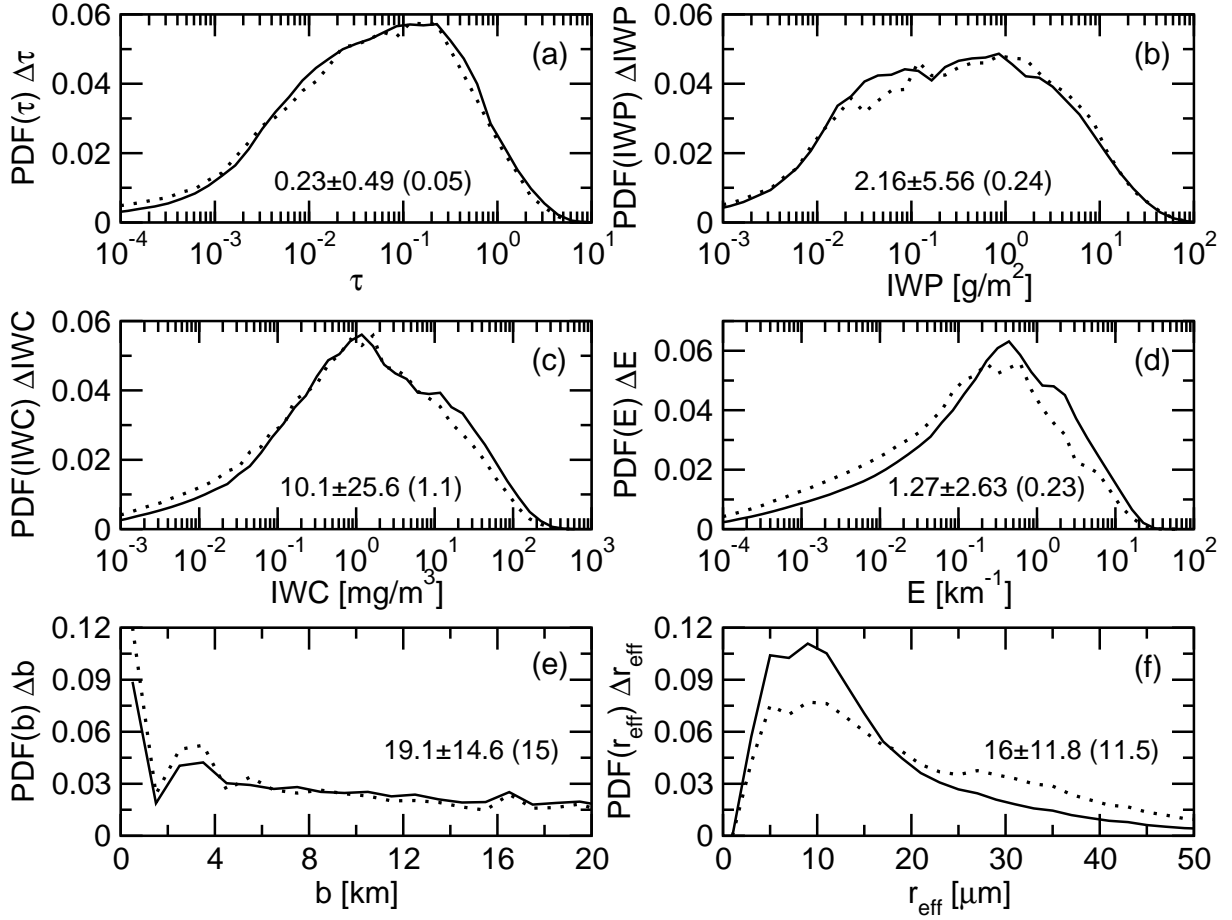


Fig. 10. Probability distributions of contrail cirrus (a) optical depth τ , (b) ice water paths IWP, (c) ice water contents IWC, (d) visible extinction E , (e) widths b , and (f) effective ice crystal radii r_{eff} . Values in each panel denote the distribution mean and standard deviation (median). The PDFs (solid curves) have been obtained by imposing simultaneous variability in temperature T , shear σ , ice supersaturation s , and supersaturated layer thickness d using the PDFs from Sect. 2.2.1. The set of dotted curves assumes a constant supersaturated layer thickness with the same mean value. Parameters are $T_m = 223 \pm 4$ K, $\sigma_m = 0.0025$ s $^{-1}$ and $c = 1.2$, $s_m = 0.1$, and $d_m = 1$ km, as suggested by the RUC analysis for 2001 over the continental U.S.A. All simulations further used uniform random sampling of initial ice particle number density n_0 (average 30 cm $^{-3}$) and Gamma distributions with initial mean ice crystal radii \bar{r}_0 (average 1 μm) according to Eqs. (17a) and (17b) with $\varphi_n = 1$.

important factors controlling optical depth. In that context, possible correlations between d and IWC need to be studied.

The set of dotted curves in Fig. 10 has been obtained by neglecting variability in the supersaturated layer thickness d , fixing d to the same mean value $d_m = 1$ km. Consistent with the results discussed in Sect. 3.1, the impact on PDF(τ) is minor. The changes in the PDFs of IWP, IWC, and b are likewise small. This does not hold for E and r_{eff} . Including variability in d markedly increases the probability to find, on average, brighter contrail cirrus composed of smaller ice crystals. This occurs at the expense of less bright clouds containing larger crystals. This happens because the PDF(d) favors the occurrence of layer depths that are smaller than the mean, from which the larger crystals quickly fall out and sublimate. Including d -variability increases E_m from 0.9 km $^{-1}$

to 1.27 km $^{-1}$, and decreases $r_{\text{eff},m}$ from 20.7 μm to 16 μm .

We reiterate the intercomparison issue of mean optical depth values (Sect. 4.2). The observed PDF(τ) from Fig. 9 (stepped lines) results in a mean $\tau_m = 0.26$. The simulated value (solid curve) is $\tau_m = 0.15$ because of the inclusion of cases with $\tau < 0.1$. Finally, when the full simulated PDF from Fig. 10a is considered that additionally includes the cases with $\tau > 1$, the mean value is $\tau_m = 0.23$.

In relating our simulated contrail cirrus distributions PDF(τ) to thin cirrus observations, we note that the distributions measured by ground-based Lidar at tropical and mid-latitude sites (Goldfarb et al., 2001; Immler and Schrems, 2002; Immler et al., 2008) show the typical generic shape as discussed in Sect. 3. This similarity builds further confidence in the present model approach. Specifically, the CC-

SIM PDF(τ), when plotted on a linear scale as in Fig. 9, and PDFs from Lidar observations of tropical thin cirrus clouds (Comstock et al., 2002, their Figure 3a) are strikingly similar. These tropical data were taken at the Department of Energy Atmospheric Radiation Measurement site on Nauru Island, deploying a Micropulse Lidar and a millimeter cloud Radar. The Lidar instrument, detecting backscatter signals, is expected to capture a substantial number of ice clouds with $\tau < 0.1$, while the Radar, based on microwave radiance measurements, misses most of the thin clouds because they are typically composed of small particles that are not detectable at those wavelengths.

6 Conclusions

In this paper we introduce an analytical model (CCSIM) to estimate contrail cirrus evolution over 4 hours in prescribed meteorological conditions. The major physical processes determining this evolution realized in the CCSIM comprise depositional growth and sedimentation of ice particles and their transport in a vertically sheared flow field. We demonstrate the usefulness of this approach for statistical analyses of contrail cirrus properties by comparing results of a representative numerical case study to Large-Eddy simulations.

We define probability distributions of the cloud-controlling factors temperature, wind shear, ice supersaturation and supersaturated layer thickness in extratropical conditions. Functional forms, mean values and variances chosen for these distributions are guided by observations. We use the distributions to drive a large number of analytical model simulations. In a first step, we examine the impact of individual cloud-controlling factors on probability distributions of contrail cirrus optical depth. In a second step, variability in contrail cirrus width and microphysical parameters is studied prescribing simultaneous variability in the cloud-controlling factors. Besides a typical baseline scenario, extreme cases leading to the generation of optically very thin or thick contrail cirrus are also investigated. A sensitivity study reveals a potential impact of aircraft wake processing of ice crystals on the temperature dependence of the optical depth distributions.

We compare CCSIM results with in situ and remote sensing measurements. Our results for optical depth, optical extinction, ice water content, ice water path, effective radius, ice crystal number concentration and width are judged reliable as they lie within the respective ranges bracketed by available measurements. We quantify for the first time that previous satellite measurements detect roughly one half of line-shaped persistent contrails; those with optical depths below 0.1 are largely missed. Because of the high frequency of occurrence of optically thin and subvisible contrails, their radiative impact may contribute significantly to the total and neglecting them may lead to an underestimation. Accounting for the poor detection efficiency of the passive remote sens-

ing method at low optical depths, we reproduce the statistic of optical depth inferred from one year of observations over the continental U.S.A. Our model provides further information on contrail cirrus properties that could not be observed.

Due to the similarity of our simulated optical depth distributions with those obtained by Lidar studies of optically thin cirrus, and since contrail cirrus and natural cirrus are affected by the same physical processes and controlling factors, we expect many of our findings summarized below to apply to thin cirrus in general.

1. Optical depth and key microphysical parameters of contrail cirrus exhibit large horizontal variability, with standard deviations exceeding mean values.
2. Neglecting vertical variability in contrail cirrus ice water content and effective ice crystal radius may not provide good estimates of column optical depths.
3. Supersaturation with respect to ice and vertical shear of the horizontal wind are key controlling factors determining the variability in contrail cirrus properties.
4. Mean contrail cirrus visible optical depth ranges between 0.05 – 0.5 depending on whether least or most favorable meteorological conditions for development prevail. Typical mean values are estimated to be 0.2 – 0.3.
5. The main effect of variability in supersaturation is to increase the probability of occurrence of optically thin contrails (optical depth < 0.1).
6. The main effect of variability in wind shear is to increase the contrail cirrus coverage by producing wider contrails, especially in deep supersaturated layers.
7. The main effect of variability in the thicknesses of ice-supersaturated layers is to produce brighter contrail cirrus, with little effect on the optical depth statistic.
8. Averaged over 4 hours of age, contrail cirrus microphysical properties differ from those of most natural cirrus in the extratropics.
9. A substantial fraction (typically 20 – 40%) of contrail cirrus is subvisible. Low optical depth is correlated with large coverage.
10. Previous passive satellite remote sensing detected only $\sim 50\%$ of line-shaped persistent contrails, strongly underestimating the occurrence frequency for optical depths $\lesssim 0.1$. The undetected fraction may not be negligible when estimating radiative forcing.
11. Variability in contrail cirrus optical depth and microphysical variables needs to be considered when comparing measurements and simulations.

12. Processes that depend non-linearly on contrail cirrus optical depth and other microphysical parameters should not be evaluated based on mean values alone.

Further research should aim at reducing the uncertainty in number and size of contrail ice particles after processing in the aircraft wake vortices to better characterize the initial stage of contrail cirrus development. Detailed cloud modeling should explore the role of secondary ice nucleation and implications for the ice mass budget and optical properties of older (> 4 h) contrail cirrus. The impact of contrail cirrus optical depth variability on radiation fluxes should be systematically studied. Future satellite data sets of line-shaped persistent contrail coverage and optical depth could be analyzed with the CCSIM deriving empirical detection efficiencies to correct observed optical depth statistics and better compare such data sets with climate model simulations. The CCSIM may also be used to support large-scale model validation.

Appendix A

Shape-preserving diffusional growth law

For particles undergoing radial growth, the size distribution f evolves according to $\partial_t f + \partial_r(\dot{r}f) = 0$. The effective single particle growth law \dot{r} from Eq. (2b) can be rewritten with the help of Eq. (2a):

$$\partial_t f + (\gamma/\bar{r}^2)\partial_r(rf) = 0, \quad (\text{A1})$$

or, using the auxiliary distribution $F = rf$,

$$\partial_t F + \dot{r}\partial_r F = 0. \quad (\text{A2})$$

Equation (A2) explicitly shows that F is constant and equal to its initial value F_0 along the characteristic curves $r(t)$ given by Eq. (2c).

Prescribing the initial Gamma distribution

$$f_0 \equiv f(r_0, t_0) = \frac{\lambda_0^{\mu+1}}{\Gamma(\mu+1)} r_0^\mu \exp(-\lambda_0 r_0), \quad (\text{A3})$$

we find

$$F_0 = r_0 f_0 = \frac{(\lambda_0 r_0)^{\mu+1}}{\Gamma(\mu+1)} \exp(-\lambda_0 r_0). \quad (\text{A4})$$

Recalling that $\lambda = (\mu+1)/\bar{r}$ and noting from Eqs. (2a) and (2c) that $\lambda r = \lambda_0 r_0$, we finally obtain the Gamma function Eq. (7b) as the solution by evaluating $F = F_0/r$ from Eq. (A4). Thus, as $\mu = \text{const.}$, the functional form (shape) of f remains unchanged during growth while the mode radius \bar{r} increases.

Appendix B

Fraction of ice crystals surviving the vortex phase

The fraction φ_n of ice crystal number surviving the adiabatic compression of an isolated, sinking aircraft vortex is based on LES studies (Unterstrasser et al., 2008, their Eq. 5.1 and Table 4). We have fitted the corresponding power law expression $\varphi_n(s, T)$ to the LES results within the ranges $T = 209 - 222$ K and $s = 0 - 0.2$.

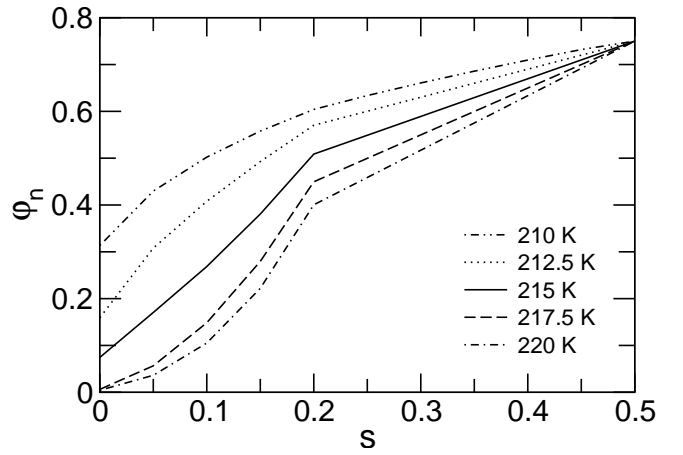


Fig. B1. Fraction of ice crystal number surviving the vortex phase versus ice supersaturation for selected temperatures.

Figure B1 shows $\varphi_n(s, T)$ as a function of s for selected temperatures used in the present work. We have extended the original formula to approximately cover a wider range of supersaturation conditions based on few further simulations. We see that φ_n is generally below unity, as some ice crystals always sublime at the vortex edges. Further, low s favors small φ_n because already small fluctuations in the saturation ratio may cause subsaturation and lead to sublimation of ice particles. Finally, φ_n decreases strongly at high T because of faster sublimation rates.

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